

**On a differential inequality for
a viscous compressible heat conducting capillary fluid
bounded by a free surface**

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Abstract. We derive a global differential inequality for solutions of a free boundary problem for a viscous compressible heat conducting capillary fluid. The inequality is essential in proving the global existence of solutions.

1. Introduction. The motion of a viscous compressible heat conducting capillary fluid in a bounded domain $\Omega_t \subset \mathbb{R}^3$ (which depends on time $t \in \mathbb{R}_+^1$) is described by the following system with the boundary and initial conditions (see [2], Chs. 2 and 5):

$$\begin{aligned}
 & \varrho[v_t + (v \cdot \nabla)v] + \nabla p - \mu\Delta v - \nu\nabla \operatorname{div} v = \varrho f && \text{in } \tilde{\Omega}^T, \\
 & \varrho_t + \operatorname{div}(\varrho v) = 0 && \text{in } \tilde{\Omega}^T, \\
 & \varrho c_v(\theta_t + v \cdot \nabla\theta) + \theta p_\theta \operatorname{div} v - \kappa\Delta\theta \\
 & \quad - \frac{\mu}{2} \sum_{i,j=1}^3 (v_{i,x_j} + v_{j,x_i})^2 - (\nu - \mu)(\operatorname{div} v)^2 = \varrho r && \text{in } \tilde{\Omega}^T, \\
 (1) \quad & \mathbb{T}\bar{n} - \sigma H\bar{n} = -p_0\bar{n} && \text{on } \tilde{S}^T, \\
 & v \cdot \bar{n} = -\phi_t/|\nabla\phi| && \text{on } \tilde{S}^T, \\
 & \partial\theta/\partial n = \theta_1 && \text{on } \tilde{S}^T, \\
 & v|_{t=0} = v_0, \quad \varrho|_{t=0} = \varrho_0, \quad \theta|_{t=0} = \theta_0 && \text{in } \Omega,
 \end{aligned}$$

where $\phi(x, t) = 0$ describes S_t , \bar{n} is the outward vector normal to the boundary (i.e. $\bar{n} = \nabla\phi/|\nabla\phi|$), $\tilde{\Omega}^T = \bigcup_{t \in (0, T)} \Omega_t \times \{t\}$, $\Omega_0 = \Omega$ is an initial domain, $\tilde{S}^T = \bigcup_{t \in (0, T)} S_t \times \{t\}$. Moreover, $v = v(x, t)$ ($v = (v_1, v_2, v_3)$) is the velocity

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of the fluid, $\varrho = \varrho(x, t)$ the density, $\theta = \theta(x, t)$ the temperature, $f = f(x, t)$ the external force field per unit mass, $r = r(x, t)$ the heat sources per unit mass, $\theta_1 = \theta_1(x, t)$ the heat flow per unit surface, $p = p(\varrho, \theta)$ the pressure, μ and ν the viscosity coefficients, κ the coefficient of heat conductivity, $c_v = c_v(\varrho, \theta)$ the specific heat at constant volume, and p_0 the external (constant) pressure. We assume that the coefficients μ, ν, κ are constants such that $\kappa > 0, \nu \geq \frac{1}{3}\mu > 0$ and moreover $c_v > 0$, which results from thermodynamic considerations. Further, $\mathbb{T} = \mathbb{T}(v, p)$ denotes the stress tensor of the form

$$\mathbb{T} = \{T_{ij}\} = \{-p\delta_{ij} + \mu(v_{i,x_j} + v_{j,x_i}) + (\nu - \mu)\delta_{ij} \operatorname{div} v\} \equiv \{-p\delta_{ij} + D_{ij}(v)\},$$

where $i, j = 1, 2, 3$, and $\mathbb{D} = \mathbb{D}(v) = \{D_{ij}\}$ is the deformation tensor.

Finally, we denote by H the double mean curvature of S_t which is negative for convex domains and can be expressed in the form

$$H\bar{n} = \Delta_{S_t}(t)x, \quad x = (x_1, x_2, x_3),$$

where $\Delta_{S_t}(t)$ is the Laplace–Beltrami operator on S_t . Let S_t be determined by $x = x(s^1, s^2, t)$, $(s^1, s^2) \in \mathbb{R}^2$. Then we have

$$\begin{aligned} \Delta_{S_t}(t) &= g^{-1/2} \frac{\partial}{\partial s^\alpha} g^{-1/2} \hat{g}_{\alpha\beta} \frac{\partial}{\partial s^\beta} \\ &= g^{-1/2} \frac{\partial}{\partial s^\alpha} g^{1/2} g^{\alpha\beta} \frac{\partial}{\partial s^\beta} \quad (\alpha, \beta = 1, 2), \end{aligned}$$

where the convention summation over repeated indices is assumed, $g = \det\{g_{\alpha\beta}\}_{\alpha, \beta=1,2}$, $g_{\alpha\beta} = x_\alpha \cdot x_\beta$, $(x^\alpha = \partial x / \partial s^\alpha)$, $\{g^{\alpha\beta}\}$ is the inverse matrix to $\{g_{\alpha\beta}\}$ and $\{\hat{g}_{\alpha\beta}\}$ is the matrix of algebraic complements of $\{g_{\alpha\beta}\}$.

Assume that the domain Ω is given. Then by (1.1)₅, $\Omega_t = \{x \in \mathbb{R}^3 : x = x(\xi, t), \xi \in \Omega\}$, where $x = x(\xi, t)$ is the solution of the Cauchy problem

$$\frac{\partial x}{\partial t} = v(x, t), \quad x|_{t=0} = \xi \in \Omega, \quad \xi = (\xi_1, \xi_2, \xi_3).$$

Hence

$$(1.2) \quad x = \xi + \int_0^t u(\xi, s) ds \equiv X_u(\xi, t),$$

where $u(\xi, t) = v(X_u(\xi, t), t)$.

Formula (1.2) yields the relation between the Eulerian x and Lagrangian ξ coordinates. Moreover, the kinematic boundary condition (1.1)₅ implies that the boundary S_t is a material surface. Thus, if $\xi \in S = S_0$ then $X_u(\xi, t) \in S_t$ and $S_t = \{x : x = X_u(\xi, t), \xi \in S\}$.

By the equation of continuity (1.1)₂ and (1.1)₅ the total mass of the drop is conserved and the following relation between ϱ and Ω_t holds:

$$\int_{\Omega_t} \varrho(x, t) dx = M.$$

In this paper we prove a global differential inequality for problem (1.1) (see Theorem 3.13) which we shall use in the next paper to prove the global-in-time existence of a solution to problem (1.1) close to a constant state. The paper is divided into three sections. In Section 2 we introduce some notation. In Section 3 we formulate a series of lemmas (see Lemmas 3.1–3.12) which are used to derive the differential inequalities (3.46) and (3.47).

Problem (1.1) is also considered in [12]–[17]. In [12] we prove the local existence of a solution to problem (1.1) in Sobolev–Slobodetskii spaces in two cases: $\sigma = 0$ and $\sigma > 0$. [14] and [15] are devoted to conservation laws for problem (1.1) in two cases: without surface tension and with it, respectively. In [14] and [15] we prove that we can choose $\varrho_0, v_0, \theta_0, \theta_1, p_0, \kappa, \sigma$ (in the case $\sigma > 0$) and the form of the internal energy per unit mass $\varepsilon = \varepsilon(\varrho, \theta)$ in such a way that $\text{var}_t |\Omega_t|$ is as small as we need. In [16] the global differential inequality in the case $\sigma = 0$, analogous to inequality (3.46) is obtained. [17] is concerned with the global-in-time existence of solutions to problem (1.1) when $\sigma = 0$. Finally, [13] contains a review of results from [12], [1]–[17] and this paper.

In order to prove the main result of the paper, i.e. Theorem 3.13, we apply the same method as in [18], [19] and [16], which is very close to the methods used in [10] and [11] (see also [4]–[7] and [8]).

Papers [18] and [19] of W. M. Zajączkowski and [9] of V. A. Solonnikov and A. Tani refer to the problem corresponding to (1.1) for a compressible barotropic fluid.

In [8] K. Pileckas and W. M. Zajączkowski proved the existence of stationary motion of a viscous compressible barotropic fluid bounded by a free surface governed by surface tension.

Finally, the motion of a viscous compressible heat conducting fluid in a fixed domain was considered by A. Matsumura and T. Nishida in [3]–[7] and by A. Valli and W. M. Zajączkowski in [11].

2. Notation. Let $Q = \Omega_t$ or $Q = S_t$ ($t \geq 0$). We denote by $\|\cdot\|_{l,Q}$ ($l \geq 0$) and $|\cdot|_{p,Q}$ ($1 \leq p \leq \infty$) the norms in the usual Sobolev spaces $W_2^l(Q)$ and $L_p(Q)$ spaces, respectively.

Next, we introduce the space $\Gamma_k^l(Q)$ of functions u with the norm

$$\|u\|_{\Gamma_k^l(Q)} = \sum_{i \leq l-k} \|\partial_t^i u\|_{l-i,Q} \equiv |u|_{l,k,Q},$$

where $l > 0$ and $k \geq 0$.

We shall use the following notation for derivatives of u . If u is a scalar-valued function we denote by $D_{x,t}^k u$ or $\underbrace{u}_{k \text{ times}}$ the vector $(D_x^\alpha \partial_t^i u)_{|\alpha|+i=k}$.

Similarly, if $u = (u_1, u_2, u_3)$ we denote by $D_{x,t}^k u$ or $\underbrace{u_{x \dots x t \dots t}}_{k \text{ times}}$ the vector $(D_x^\alpha \partial_t^i u_j)_{|\alpha|+i=k, j=1,2,3}$. Hence

$$|D_{x,t}^k u| = \sum_{|\alpha|+i=k} |D_x^\alpha \partial_t^i u|.$$

We use the following lemma:

LEMMA 2.1. *The following imbedding holds:*

$$W_r^l(Q) \subset L_p^\alpha(Q) \quad (Q \subset \mathbb{R}^3),$$

where $|\alpha| + 3/r - 3/p \leq l$, $l \in \mathbb{Z}$, $1 \leq p, r \leq \infty$; $L_p^\alpha(Q)$ is the space of functions u such that $|D_x^\alpha u|_{p,Q} < \infty$, and $W_r^l(Q)$ is the Sobolev space.

Moreover, the following interpolation inequalities are true:

$$(2.1) \quad |D_x^\alpha u|_{p,Q} \leq c\varepsilon^{1-\kappa} |D_x^l u|_{r,Q} + c\varepsilon^{-\kappa} |u|_{r,Q},$$

where $\kappa = |\alpha|/l + 3/(lr) - 3/(lp) < 1$, ε is a parameter, $c > 0$ is a constant independent of u and ε ;

$$(2.2) \quad |D_x^\alpha u|_{q,\partial Q} \leq c\varepsilon^{1-\kappa} |D_x^l u|_{r,Q} + c\varepsilon^{1-\kappa} |u|_{r,Q},$$

where $\kappa = |\alpha|/l + 3/(lr) - 2/(lq) < 1$, ε is a parameter, $c > 0$ is a constant independent of u and ε . ■

Lemma 2.1 follows from Theorem 10.2 of [1].

3. Global differential inequality. In this section we assume that $\nu > \frac{1}{3}\mu$. Further, assume that the existence of a sufficiently smooth local solution of problem (1.1) has been proved. To prove the desired differential inequality we assume that Ω_t ($t \leq T$, T is the time of local existence) is diffeomorphic to a ball, so S_t can be described by

$$|x| \equiv r = R(\omega, t), \quad \omega \in S^1,$$

where S^1 is the unit sphere.

Moreover, we consider the motion near the constant state $v_e = 0$, $p_e = p_0 + 2\sigma/R_e$, $\theta_e = (1/|\Omega|) \int_\Omega \theta_0 d\xi$, $\varrho_e = M/((4\pi/3)R_e^3)$, where R_e is a solution of the equation

$$p\left(\frac{M}{(4\pi/3)R_e^3}, \theta_e\right) = p_e.$$

(Obviously, we assume that the above equation is solvable with respect to $R_e > 0$.)

Let

$$p_\sigma = p - p_0 - q_0, \quad \varrho_\sigma = \varrho - \varrho_e, \quad \vartheta_0 = \theta - \theta_e, \quad \vartheta = \theta - \theta_{\Omega_t},$$

where $q_0 = 2\sigma/R_e$ and $\theta_{\Omega_t} = (1/|\Omega_t|) \int_{\Omega_t} \theta dx$. Then problem (1.1) takes the form

$$(3.1) \quad \begin{aligned} & \varrho[v_t + (v \cdot \nabla)v] - \operatorname{div} \mathbb{T}(v, p_\sigma) = \varrho f && \text{in } \Omega_t, \quad t \in [0, T], \\ & \varrho_t + \operatorname{div}(\varrho v) = 0 && \text{in } \Omega_t, \quad t \in [0, T], \\ & \varrho c_v(\varrho, \theta)(\vartheta_{0t} + v \cdot \nabla \vartheta_0) + \theta p_\theta(\varrho, \theta) \operatorname{div} v \\ & \quad - \kappa \Delta \vartheta_0 - \frac{\mu}{2} \sum_{i,j} (\partial_{x_i} v_j + \partial_{x_j} v_i)^2 \\ & \quad - (\nu - \mu)(\operatorname{div} v)^2 = \varrho r && \text{in } \Omega_t, \quad t \in [0, T], \\ & \mathbb{T}(v, p_\sigma) \bar{n} = \sigma \Delta_{S_t} x \cdot \bar{n} \bar{n} + q_0 \bar{n} && \text{on } S_t, \quad t \in [0, T], \\ & \partial \vartheta_0 / \partial n = \theta_1 && \text{on } S_t, \quad t \in [0, T], \end{aligned}$$

where $\mathbb{T}(v, p_\sigma) = \{\mu(\partial_{x_i} v_j + \partial_{x_j} v_i) + (\nu - \mu)\delta_{ij} \operatorname{div} v - p_\sigma \delta_{ij}\}$ and T is the time of local existence.

In the sequel we shall use the following Taylor formula for p_σ :

$$(3.2) \quad \begin{aligned} p_\sigma &= p(\varrho, \theta) - p(\varrho_e, \theta_e) \\ &= p(\varrho, \theta) - p(\varrho_e, \theta) + p(\varrho_e, \theta) - p(\varrho_e, \theta_e) \\ &= (\varrho - \varrho_e) \int_0^1 p_\varrho(\varrho_e + s(\varrho - \varrho_e), \theta) ds \\ &\quad + (\theta - \theta_e) \int_0^1 p_\theta(\varrho_e, \theta_e + s(\theta - \theta_e)) ds \\ &\equiv p_1 \varrho_\sigma + p_2 \vartheta_0. \end{aligned}$$

We shall also use the formula:

$$(3.3) \quad \begin{aligned} p_\sigma &= p(\varrho, \theta) - p(\varrho_{\Omega_t}, \theta_{\Omega_t}) \\ &= (\varrho - \varrho_{\Omega_t}) \int_0^1 p_\varrho(\varrho_{\Omega_t} + s(\varrho - \varrho_{\Omega_t}), \theta) ds \\ &\quad + (\theta - \theta_{\Omega_t}) \int_0^1 p_\theta(\varrho_{\Omega_t}, \theta_{\Omega_t} + s(\theta - \theta_{\Omega_t})) ds \\ &\equiv p_3 \bar{\varrho}_{\Omega_t} + p_4 \vartheta, \end{aligned}$$

where the function $\varrho_{\Omega_t} = \varrho_{\Omega_t}(t)$ is a solution of the problem

$$(3.4) \quad p(\varrho_{\Omega_t}, \theta_{\Omega_t}) = p_e, \quad \varrho_{\Omega_t}|_{t=0} = \varrho_e$$

and

$$\bar{\varrho}_{\Omega_t} = \varrho - \varrho_{\Omega_t}.$$

The functions p_i ($i = 1, 2, 3, 4$) in (3.4) and (3.5) are positive.

Set

$$\begin{aligned}\varrho_* &= \min_{\tilde{\Omega}^T} \varrho(x, t), & \varrho^* &= \max_{\tilde{\Omega}^T} \varrho(x, t), \\ \theta_* &= \min_{\tilde{\Omega}^T} \theta(x, t), & \theta^* &= \max_{\tilde{\Omega}^T} \theta(x, t).\end{aligned}$$

Now we point out the following facts concerning the estimates in Lemmas 3.1–3.12 and Theorem 3.13:

1. We denote by ε small constants and for simplicity we do not distinguish them.

2. We denote by C_1 and C_2 constants which depend on ϱ_* , ϱ^* , θ_* , θ^* , T , $\int_0^T \|v\|_{3,\Omega_t}^2 dt'$, $\|S\|_{4+1/2}$, on the parameters which guarantee the existence of the inverse transformation to $x = x(\xi, t)$ and also on the constants of imbedding theorems and Korn inequalities. C_1 is always the coefficient of a linear term, while C_2 is the coefficient of a nonlinear term. For simplicity we do not distinguish different C_1 's and C_2 's.

3. We denote by c absolute constants which may depend on such parameters as μ , ν , κ , and by $c_0 < 1$ positive constants which may depend on $\mu, \nu, \kappa, \varrho_*, \varrho^*, \theta_*, \theta^*$. For simplicity we do not distinguish different C 's and C_0 's.

4. We underline that all the estimates are obtained under the assumption that there exists a local-in-time solution of (1.1), so all the quantities ϱ_* , ϱ^* , θ_* , θ^* , T , $\int_0^T \|v\|_{3,\Omega_t}^2 dt'$, $\|S\|_{4+1/2}$ are estimated by the data functions. Moreover, the existence of the inverse transformation to $x = x(\xi, t)$ is guaranteed by the estimates for the local solutions (see [12]).

LEMMA 3.1. *Let v, ϱ, ϑ_0 be a sufficiently smooth solution of (3.1). Then*

$$\begin{aligned}(3.6) \quad & \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v^2 + \frac{p_1}{\varrho} \varrho_\sigma^2 + \bar{\varrho}_{\Omega_t}^2 + \frac{p_2 \varrho c_v}{p_\theta \theta} \vartheta_0^2 \right) dx \\ & + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \left(g^{\alpha\beta} \bar{n} \cdot \int_0^t v_{s^\alpha} dt' \bar{n} \cdot \int_0^t v_{s^\beta} dt' \right) ds + c_0 \|v\|_{1,\Omega_t}^2 \\ & + (\nu - \mu) \|\operatorname{div} v\|_{0,\Omega_t}^2 + c_0 \|\vartheta_{0x}\|_{0,\Omega_t}^2 \\ & \leq \varepsilon \left(\|p_\sigma\|_{0,\Omega_t}^2 + \|\vartheta_{0tx}\|_{0,\Omega_t}^2 + \left\| \int_0^t v_s dt' \right\|_{0,S_t}^2 + \|H(\cdot, 0) + 2/R_e\|_{0,S^1}^2 \right) \\ & + C_1 (\|v\|_{0,\Omega_t}^2 + \|r\|_{0,\Omega_t}^2 + \|r\|_{0,\Omega_t} + \|\theta_1\|_{1,\Omega_t}^2 \\ & + \|\theta_1\|_{1,\Omega_t} + \|f\|_{0,\Omega_t}^2) + C_2 X_1 Y_1,\end{aligned}$$

where $\varepsilon > 0$ is sufficiently small, $s = (s^1, s^2)$ and

$$\begin{aligned} X_1 &= \|v\|_{2,\Omega_t}^2 + \|\varrho_\sigma\|_{2,\Omega_t}^2 + \|\vartheta_0\|_{2,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2, \\ Y_1 &= X_1 + \left\| \int_0^1 v dt' \right\|_{2,S_t}^2. \end{aligned}$$

P r o o f. Multiplying (3.1)₁ by v , integrating over Ω_t and using the continuity equation (3.1)₂, boundary condition (3.1)₄ and (3.2) we obtain

$$\begin{aligned} (3.7) \quad & \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \varrho v^2 dx + \frac{\mu}{2} E_{\Omega_t}(v) + (\nu - \mu) \|\operatorname{div} v\|_{0,\Omega_t}^2 \\ & - \int_{\Omega_t} p_1 \varrho_\sigma \operatorname{div} v dx - \int_{\Omega_t} p_2 \vartheta_0 \operatorname{div} v dx \\ & - \sigma \int_{S_t} (\Delta_{S_t} x \cdot \bar{n} + 2/R_e) \bar{n} \cdot v ds = \int_{\Omega_t} \varrho f v dx, \end{aligned}$$

where $E_{\Omega_t}(v) = \int_{\Omega_t} \sum_{i,j=1}^3 (\partial_{x_i} v_j + \partial_{x_j} v_i)^2 dx$.

First, we consider the sum of the second and third terms on the left-hand side of (3.7). We have

$$\begin{aligned} (3.8) \quad & \frac{\mu}{2} E_{\Omega_t}(v) + (\nu - \mu) \|\operatorname{div} v\|_{0,\Omega_t}^2 \\ & = \frac{\mu}{2} \int_{\Omega_t} (v_{i,x_j} + v_{j,x_i})^2 dx + (\nu - \mu) \int_{\Omega_t} (\operatorname{div} v)^2 dx \\ & = \frac{\mu}{2} \sum_{i \neq j} \int_{\Omega_t} (v_{i,x_j} + v_{j,x_i})^2 dx + \frac{\mu}{2} \sum_{i=j} \int_{\Omega_t} (v_{i,x_j} + v_{j,x_i})^2 dx \\ & \quad + (\nu - \mu) \int_{\Omega_t} (\operatorname{div} v)^2 dx \\ & = \frac{\mu}{2} \sum_{i \neq j} \int_{\Omega_t} (v_{i,x_j} + v_{j,x_i})^2 dx + \frac{\mu}{2} \varepsilon_1 \sum_{i=j} \int_{\Omega_t} (v_{i,x_j} + v_{j,x_i})^2 dx \\ & \quad + \frac{\mu}{2} (1 - \varepsilon_1) \cdot 4 \sum_i \int_{\Omega_t} (v_{i,x_j})^2 dx + (\nu - \mu) \int_{\Omega_t} (\operatorname{div} v)^2 dx \equiv I, \end{aligned}$$

where $\varepsilon_1 \in (0, 1)$. Since $(\xi_1 + \xi_2 + \xi_3)^2 \leq 3(\xi_1^2 + \xi_2^2 + \xi_3^2)$ the last two terms in I are estimated from below by

$$[\nu - (1 + 2\varepsilon_1)\mu/2] \int_{\Omega_t} (\operatorname{div} v)^2 dx.$$

Assuming that $\nu = (1 + 2\varepsilon_1)\mu/3$ we obtain $\varepsilon_1 = \frac{3}{2\mu}(\nu - \mu/3)$, so

$$(3.9) \quad I \geq \frac{\mu}{2}\varepsilon_1 \int_{\Omega_t} (v_{i,x_j} + v_{j,x_i})^2 dx = \frac{3}{4} \left(\nu - \frac{\mu}{3} \right) \int_{\Omega_t} (v_{i,x_j} + v_{j,x_i})^2 dx.$$

By the continuity equation (3.1)₂, energy equation (3.1)₃ and boundary condition (3.1)₅ we have

$$(3.10) \quad \begin{aligned} - \int_{\Omega_t} p_1 \varrho_\sigma \operatorname{div} v dx &= \int_{\Omega_t} \frac{p_1}{\varrho} \varrho_\sigma (\varrho_{\sigma_t} + v \cdot \nabla \varrho_\sigma) dx \\ &= \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \frac{p_1 \varrho_\sigma^2}{\varrho} dx + I_1, \end{aligned}$$

where

$$(3.11) \quad \begin{aligned} |I_1| &\leq \varepsilon (\|v_x\|_{0,\Omega_t}^2 + \|\vartheta_{0x}\|_{0,\Omega_t}^2) + C_1 (\|r\|_{0,\Omega_t}^2 + \|\theta_1\|_{1,\Omega_t}^2) \\ &\quad + C_2 (\|\varrho_\sigma\|_{1,\Omega_t}^4 + \|v\|_{1,\Omega_t}^2 \|\varrho_\sigma\|_{2,\Omega_t}^2 + \|v\|_{1,\Omega_t}^2 \|\vartheta_0\|_{2,\Omega_t}^2 \\ &\quad + \|v\|_{2,\Omega_t}^2 \|\varrho_\sigma\|_{1,\Omega_t}^2 + \|\varrho_\sigma\|_{2,\Omega_t}^2 \|\varrho_\sigma\|_{1,\Omega_t}^2). \end{aligned}$$

Now, we consider the boundary term in (3.7). In the same way as in Lemma 4.1 of [19] we obtain

$$(3.12) \quad \begin{aligned} - \int_{S_t} (\Delta_{S_t} x \cdot \bar{n} + 2/R_e) v \cdot \bar{n} ds \\ &= \frac{1}{2} \frac{d}{dt} \int_{S_t} \left(g^{\alpha\beta} \bar{n} \cdot \int_0^t v_{s^\alpha} dt' \bar{n} \cdot \int_0^t v_{s^\beta} dt' \right) ds + I_1, \end{aligned}$$

where

$$(3.13) \quad \begin{aligned} |I_1| &\leq \varepsilon \left(\left\| \int_0^t v_s dt' \right\|_{0,S_t}^2 + \|H(\cdot, 0) + 2/R_e\|_{0,S^1}^2 + \|v\|_{1,\Omega_t}^2 \right) \\ &\quad + C_1 \|v\|_{0,\Omega_t}^2 + C_2 \left\| \int_0^t v dt' \right\|_{2,S_t}^2 \|v\|_{2,\Omega_t}^2. \end{aligned}$$

Next, dividing (3.1)₃ by θp_θ , multiplying the result by $p_2 \vartheta_0$ and integrating over Ω_t we get

$$(3.14) \quad \begin{aligned} &\int_{\Omega_t} \frac{p_2 \varrho c_v}{\theta p_\theta} \left(\partial_t \frac{\vartheta_0^2}{2} + v \cdot \nabla \frac{\vartheta_0^2}{2} \right) dx + \int_{\Omega_t} p_2 \vartheta_0 \operatorname{div} v dx \\ &- \int_{\Omega_t} \frac{p_2 \kappa \Delta \vartheta_0}{\theta p_\theta} \vartheta_0 dx - \int_{\Omega_t} \frac{p_2 \mu}{2 \theta p_\theta} \sum_{i,j} (\partial_{x_j} v_i + \partial_{x_i} v_j)^2 \vartheta_0 dx \\ &- \int_{\Omega_t} \frac{p_2 (\nu - \mu)}{\theta p_\theta} (\operatorname{div} v)^2 \vartheta_0 dx = \int_{\Omega_t} \frac{p_2 \varrho r}{\theta p_\theta} \vartheta_0 dx. \end{aligned}$$

Therefore, using the same argument as in Lemma 3.1 of [16], by (3.7)–(3.14) we have

$$\begin{aligned}
(3.15) \quad & \frac{1}{2} \frac{d}{dt} \left(\varrho v^2 + \frac{p_1 \varrho_\sigma^2}{\varrho} + \frac{p_2 \varrho c_v}{\theta p_\theta} \vartheta_0^2 \right) dx \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t}^t g^{\alpha\beta} \bar{n} \cdot \int_0^t v_{s^\alpha} dt' \bar{n} \cdot \int_0^t v_{s^\beta} dt' ds + c_0 \|v\|_{1,\Omega_t}^2 \\
& + (\nu - \mu) \|\operatorname{div} v\|_{0,\Omega_t}^2 + c_0 \|\vartheta_{0x}\|_{0,\Omega_t}^2 \\
& \leq \varepsilon \left(\left\| \int_0^t v_s dt' \right\|_{0,S_t}^2 + \|H(\cdot, 0) + 2/R_0\|_{0,S^1}^2 \right) \\
& + C_1 (\|v\|_{0,\Omega_t}^2 + \|r\|_{0,\Omega_t}^2 + \|r\|_{0,\Omega_t} \\
& + \|\theta_1\|_{1,\Omega_t}^2 + \|\theta_1\|_{1,\Omega_t} + \|f\|_{0,\Omega_t}^2) + C_2 X_1 Y_1.
\end{aligned}$$

Finally, by (3.1)₂ and (3.5) we have

$$(3.16) \quad \partial_t \bar{\varrho}_{\Omega_t} + v \cdot \nabla \bar{\varrho}_{\Omega_t} + \varrho \operatorname{div} v + \partial_t \varrho_{\Omega_t} = 0,$$

where in view of (3.4) we get

$$(3.17) \quad \partial_t \varrho_{\Omega_t} = - \frac{p_{\theta_{\Omega_t}}}{p_{\varrho_{\Omega_t}}} \partial_t \theta_{\Omega_t}.$$

Using the definition of θ_{Ω_t} we calculate

$$\begin{aligned}
(3.18) \quad & \partial_t \theta_{\Omega_t} = \frac{1}{|\Omega_t|} \int_{\Omega_t} \vartheta_{0t} dx + \frac{1}{|\Omega_t|} \int_{\Omega_t} \theta \operatorname{div} v dx \\
& - \frac{1}{|\Omega_t|^2} \left(\int_{\Omega_t} \theta dx \right) \left(\int_{\Omega_t} \operatorname{div} v dx \right).
\end{aligned}$$

Equation (3.16) and formulas (3.17), (3.18) yield

$$\begin{aligned}
(3.19) \quad & \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \bar{\varrho}_{\Omega_t}^2 dx \leq \varepsilon (\|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\vartheta_{0tx}\|_{0,\Omega_t}^2) \\
& + C_1 (\|v_x\|_{0,\Omega_t}^2 + \|\vartheta_{0x}\|_{0,\Omega_t}^2 + \|r\|_{0,\Omega_t}^2 + \|\theta_1\|_{1,\Omega_t}^2) \\
& + C_2 (\|v\|_{1,\Omega_t}^2 \|\vartheta_0\|_{2,\Omega_t}^2 + \|v\|_{2,\Omega_t}^4 + \|\bar{\varrho}_{\Omega_t}\|_{1,\Omega_t}^4 \\
& + \|\varrho_\sigma\|_{2,\Omega_t}^2 \|\vartheta_0\|_{2,\Omega_t}^2 + \|\vartheta_0\|_{2,\Omega_t}^4).
\end{aligned}$$

By (3.3) and the Poincaré inequality

$$(3.20) \quad \|\vartheta\|_{0,\Omega_t} \leq \|\vartheta_{0x}\|_{0,\Omega_t}$$

we get

$$(3.21) \quad \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t} \leq C_1 (\|\vartheta_{0x}\|_{0,\Omega_t} + \|p_\sigma\|_{0,\Omega_t}).$$

The estimates (3.15), (3.19) and (3.21) yield (3.6). ■

LEMMA 3.2. Let v, ϱ, ϑ_0 be a sufficiently smooth solution of (3.1). Then

$$\begin{aligned}
(3.22) \quad & \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_t^2 + \frac{p_{\sigma\varrho}}{\varrho} \varrho_{\sigma t}^2 + \frac{\varrho c_v}{\theta} \vartheta_{0t}^2 \right) dx \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} g^{\alpha\beta} v_{s^\alpha} \cdot \bar{n} v_{s^\beta} \cdot \bar{n} ds + c_0 \|v_t\|_{1,\Omega_t}^2 \\
& + (\nu - \mu) \|\operatorname{div} v_t\|_{0,\Omega_t}^2 + c_0 \|\vartheta_{0t}\|_{1,\Omega_t}^2 \\
& \leq \varepsilon (\|v\|_{1,\Omega_t}^2 + \|v_{xx}\|_{0,\Omega_t}^2 + \|\vartheta_{0x}\|_{0,\Omega_t}^2) \\
& + C_1 (|f|_{1,0,\Omega_t}^2 + |r|_{1,0,\Omega_t}^2 + \|r\|_{0,\Omega_t} + |\theta_1|_{2,1,\Omega_t}^2 + \|\theta_1\|_{1,\Omega_t}) \\
& + C_2 X_2^2 (1 + X_2),
\end{aligned}$$

where $X_2 = |v|_{2,1,\Omega_t}^2 + |\varrho_\sigma|_{2,1,\Omega_t}^2 + |\vartheta_0|_{2,1,\Omega_t}^2$.

Proof. By the same argument as in Lemma 3.2 of [16] we have

$$\begin{aligned}
(3.23) \quad & \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_t^2 + \frac{p_{\sigma\varrho}}{\varrho} \varrho_{\sigma t}^2 + \frac{\varrho c_v}{\theta} \vartheta_{0t}^2 \right) dx + \|v_t\|_{1,\Omega_t}^2 \\
& + (\nu - \mu) \|\operatorname{div} v_t\|_{0,\Omega_t}^2 + \|\vartheta_{0t}\|_{1,\Omega_t}^2 - \int_{S_t} [\mathbb{T}(v, p_\sigma)]_{,t} \bar{n} \cdot v_t ds \\
& \leq \varepsilon (\|\varrho_{\sigma t}\|_{0,\Omega_t}^2 + \|v_t\|_{0,\Omega_t}^2 + \|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|\vartheta_{0tx}\|_{0,\Omega_t}^2) \\
& + C_1 (\|r\|_{0,\Omega_t}^2 + \|r_t\|_{0,\Omega_t}^2 + |\theta_1|_{2,1,\Omega_t}^2) + C_2 X_2^2 (1 + X_2).
\end{aligned}$$

By the boundary condition (3.1)₄ and the same argument as in Lemma 4.2 of [19] we get

$$(3.24) \quad - \int_{S_t} [\mathbb{T}(v, p_\sigma)]_{,t} \bar{n} \cdot v_t ds = \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} g^{\alpha\beta} v_{s^\alpha} \cdot \bar{n} v_{s^\beta} \cdot \bar{n} ds + I_2,$$

where

$$(3.25) \quad |I_2| \leq \varepsilon (\|v_t\|_{1,\Omega_t}^2 + \|v_{xx}\|_{0,\Omega_t}^2) + C_1 \|v\|_{0,\Omega_t}^2 + C_2 X_2^2.$$

From the continuity equation (3.1)₂ it follows that

$$(3.26) \quad \|\varrho_{\sigma t}\|_{0,\Omega_t}^2 \leq C_1 \|v\|_{1,\Omega_t}^2 + C_2 \|v\|_{1,\Omega_t}^2 \|\varrho_\sigma\|_{2,\Omega_t}^2$$

and equation (3.1)₃ yields

$$\begin{aligned}
(3.27) \quad & \|\vartheta_{0t}\|_{0,\Omega_t}^2 \leq \varepsilon \|\vartheta_{0xt}\|_{0,\Omega_t}^2 \\
& + C_1 (\|v_x\|_{0,\Omega_t}^2 + \|\vartheta_{0x}\|_{0,\Omega_t}^2 + \|r\|_{0,\Omega_t}^2 + |\theta_1|_{1,\Omega_t}^2) \\
& + C_2 (\|v\|_{1,\Omega_t}^2 \|\vartheta_0\|_{2,\Omega_t}^2 + \|v\|_{1,\Omega_t}^4 \\
& + \|\varrho_\sigma\|_{2,\Omega_t}^2 \|\vartheta_0\|_{2,\Omega_t}^2 + \|\vartheta_0\|_{2,\Omega_t}^4).
\end{aligned}$$

Therefore, taking into account (3.23)–(3.27) we obtain (3.22). ■

Lemmas 3.1 and 3.2 imply

LEMMA 3.3. *Let v, ϱ, ϑ_0 be a sufficiently smooth solution of (3.1). Then*

$$\begin{aligned}
& \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left[\varrho(v^2 + v_t^2) + \frac{1}{\varrho} (p_1 \varrho_\sigma^2 + p_{\sigma\varrho} \varrho_{\sigma t}^2) + \bar{\varrho}_{\Omega_t}^2 + \frac{\varrho c_v}{\theta} \left(\frac{p_2}{p_\theta} \vartheta_0^2 + \vartheta_{0t}^2 \right) \right] dx \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} g^{\alpha\beta} \left[\bar{n} \cdot \int_0^t v_{s^\alpha} dt' \bar{n} \cdot \int_0^t v_{s^\beta} dt' + \bar{n} \cdot v_{s^\alpha} \bar{n} \cdot v_{s^\beta} \right] ds \\
& + c_0 (\|v\|_{1,\Omega_t}^2 + \|v_t\|_{1,\Omega_t}^2) + (\nu - \mu) (\|\operatorname{div} v\|_{0,\Omega_t}^2 + \|\operatorname{div} v_t\|_{0,\Omega_t}^2) \\
& + c_0 (\|\vartheta_{0x}\|_{0,\Omega_t}^2 + \|\vartheta_{0t}\|_{1,\Omega_t}^2) \\
& \leq \varepsilon \left(\|\varrho_\sigma\|_{0,\Omega_t}^2 + \|v_{xx}\|_{0,\Omega_t}^2 + \left\| \int_0^t v_s dt' \right\|_{0,S_t}^2 + \|H(\cdot, 0) + 2/R_e\|_{0,S^1}^2 \right) \\
& + C_1 (\|v\|_{0,\Omega_t}^2 + |r|_{1,0,\Omega_t}^2 + \|r\|_{0,\Omega_t} + |\theta_1|_{2,1,\Omega_t}^2 + \|\theta_1\|_{1,\Omega_t} + |f|_{1,0,\Omega_t}^2) \\
& + C_2 [X_1 Y_1 + X_2^2 (1 + X_2)]. \blacksquare
\end{aligned}$$

In order to obtain an inequality for derivatives with respect to x we rewrite problem (3.1) in the Lagrangian coordinates and next we introduce a partition of unity in the fixed domain Ω . Thus we have

$$\begin{aligned}
(3.28) \quad & \eta u_{it} - \nabla_{u_j} \mathbb{T}_u^{ij}(u, p_\sigma) = \eta g_i, \quad i = 1, 2, 3, \\
& \eta_{\sigma_t} + \eta \nabla_u \cdot u = 0, \\
& \eta c_v(\eta, \Gamma) \gamma_{0t} - \kappa \nabla_u^2 \gamma_0 = \eta k - \Gamma p_\Gamma(\eta, \Gamma) \nabla_u \cdot u \\
& + \frac{\mu}{2} \sum_{i,j=1}^3 (\xi_{kx_i} \partial_{\xi_k} u_j + \xi_{kx_j} \partial_{\xi_k} u_i)^2 \\
& + (\nu - \mu) (\nabla_u \cdot u)^2,
\end{aligned}$$

$$\mathbb{T}_u(u, p_\sigma) \bar{n} = \sigma \Delta_{S_t} x(\xi, t) \cdot \bar{n} \bar{n} + q_0 \bar{n},$$

$$\bar{n} \cdot \nabla_u \gamma_0 = \Gamma_1,$$

where $\eta(\xi, t) = \varrho(x(\xi, t), t)$, $u(\xi, t) = v(x(\xi, t), t)$, $g(\xi, t) = f(x(\xi, t), t)$, $\Gamma(\xi, t) = \theta(x(\xi, t), t)$, $\gamma_0(\xi, t) = \vartheta_0(x(\xi, t), t)$, $\Gamma_1(\xi, t) = \theta_1(x(\xi, t), t)$, $\bar{n} = \bar{n}(\xi, t)$ and

$$\mathbb{T}_u(u, p_\sigma) = \{T_u^{ij}(u, p_\sigma)\} = \{-p_\sigma \delta_{ij} + \mu(\nabla_{u_i} u_j + \nabla_{u_j} u_i) + (\nu - \mu) \delta_{ij} \nabla_u \cdot u\},$$

$$\nabla_{u_i} = \xi_{kx_i} \partial_{\xi_k} \text{ and } \operatorname{div} \mathbb{T}_u(u, p_\sigma) = \nabla_u \cdot \mathbb{T}_u(u, p_\sigma).$$

By (3.4) and (3.5) we have respectively

$$p_\sigma = p_1 \eta_\sigma + p_2 \gamma_0$$

and

$$p_\sigma = p_3 \bar{\eta}_{\Omega_t} + p_4 \gamma,$$

where $\eta_\sigma = \eta - \varrho_e$, $\gamma_0 = \Gamma - \theta_e$, $\bar{\eta}_{\Omega_t} = \eta - \varrho_{\Omega_t}$, $p_1 = \int_0^1 p_\eta(\varrho_e + s\eta_\sigma, \Gamma) ds$, $p_2 = \int_0^1 p_\Gamma(\varrho_e, \theta_e + s\gamma_0) ds$, $p_3 = \int_0^1 p_\eta(\varrho_{\Omega_t} + s\bar{\eta}_{\Omega_t}, \Gamma) ds$, $p_4 = \int_0^1 p_\Gamma(\varrho_{\Omega_t}, \theta_{\Omega_t} + s\gamma) ds$, $p_i > 0$ ($i = 1, 2, 3, 4$).

Let us introduce a partition of unity $(\{\tilde{\Omega}_i\}, \{\zeta_i\})$, $\Omega = \bigcup_i \tilde{\Omega}_i$. Let $\tilde{\Omega}$ be one of the $\tilde{\Omega}_{i,s}$ and $\zeta(\xi) = \zeta_i(\xi)$ be the corresponding function. If $\tilde{\Omega}$ is an interior subdomain then let $\tilde{\omega}$ be a set such that $\tilde{\omega} \subset \tilde{\Omega}$ and $\zeta(\xi) = 1$ for $\xi \in \tilde{\omega}$. Otherwise we assume that $\overline{\tilde{\Omega}} \cap S = \emptyset$, $\tilde{\omega} \cap S \neq \emptyset$, $\tilde{\omega} \subset \overline{\tilde{\Omega}}$. Take any $\beta \in \tilde{\omega} \cap S \subset \overline{\tilde{\Omega}} \cap S = \tilde{S}\partial$ and introduce local coordinates $\{y\}$ associated with $\{\xi\}$ by the relation

$$(3.29) \quad y_k = \alpha_{kl}(\xi_l - \beta_l), \quad \alpha_{3k} = n_k(\beta), \quad k = 1, 2, 3,$$

where $\{\alpha_{kl}\}$ is a constant orthogonal matrix such that \tilde{S} is described by the equation $y_3 = F(y_1, y_2)$, $F \in W_2^{4-1/2}$ and

$$\tilde{\Omega} = \{y : |y_i| < d, i = 1, 2, F(y') < y_3 < F(y') + d, y' = (y_1, y_2)\}.$$

Next introduce functions u' , η' , Γ' , γ'_0 , γ' , Γ'_1 by means of the formulas

$$\begin{aligned} u'_i(y) &= \alpha_{ij} u_j(\xi)|_{\xi=\xi(y)}, & \eta'(y) &= \eta(\xi)|_{\xi=\xi(y)}, \\ \Gamma'(y) &= \Gamma(\xi)|_{\xi=\xi(y)}, & \gamma'_0(y) &= \gamma_0(\xi)|_{\xi=\xi(y)}, \\ \gamma'(y) &= \gamma(\xi)|_{\xi=\xi(y)}, & \Gamma'_1(y) &= \Gamma_1(\xi)|_{\xi=\xi(y)}, \end{aligned}$$

where $\xi = \xi(y)$ is the inverse transformation to (3.29). Further, we introduce new variables by

$$z_i = y_i \quad (i = 1, 2), \quad z_3 = y_3 - \tilde{F}(y), \quad y \in \tilde{\Omega},$$

which will be denoted by $z = \Phi(y)$, where \tilde{F} is an extension of F , so $\tilde{F} \in W_2^4$.

Let $\hat{\Omega} = \Phi(\tilde{\Omega}) = \{z : |z_i| < d, i = 1, 2, 0 < z_3 < d\}$ and $\hat{S} = \Phi(S)$. Define

$$\begin{aligned} \hat{u}(z) &= u'(y)|_{y=\Phi^{-1}(z)}, & \hat{\eta}(z) &= \eta'(y)|_{y=\Phi^{-1}(z)}, \\ \hat{\Gamma}(z) &= \Gamma'(y)|_{y=\Phi^{-1}(z)}, & \hat{\gamma}_0(z) &= \gamma'_0(y)|_{y=\Phi^{-1}(z)}, \\ \hat{\gamma}(z) &= \gamma'(y)|_{y=\Phi^{-1}(z)}, & \hat{\Gamma}_1(z) &= \Gamma_1(y)|_{y=\Phi^{-1}(z)}. \end{aligned}$$

Set $\hat{\nabla}_k = \xi_{lx_k}(\xi) z_i \xi_l \nabla_{z_i}|_{\xi=\chi^{-1}(z)}$, where $\chi(\xi) = \Phi(\psi(\xi))$ and $y = \psi(\xi)$ is described by (3.29). We also introduce the following notation:

$$\begin{aligned} \tilde{u}(\xi) &= u(\xi)\zeta(\xi), & \tilde{\eta}(\xi) &= \eta(\xi)\zeta(\xi), \\ \tilde{\Gamma}(\xi) &= \Gamma(\xi)\zeta(\xi), & \tilde{\gamma}_0(\xi) &= \gamma_0(\xi)\zeta(\xi), \\ \tilde{\gamma}(\xi) &= \gamma(\xi)\zeta(\xi), & \tilde{\Gamma}_1(\xi) &= \Gamma_1(\xi)\zeta(\xi) \end{aligned}$$

for $\xi \in \widetilde{\Omega}$, $\widetilde{\Omega} \cap S = \emptyset$ and

$$\begin{aligned}\tilde{u}(z) &= \widehat{u}(z)\widehat{\zeta}(z), & \tilde{\eta}(z) &= \widehat{\eta}(z)\widehat{\zeta}(z), \\ \widetilde{\Gamma}(z) &= \widehat{\Gamma}(z)\widehat{\zeta}(z), & \widetilde{\gamma}_0(z) &= \widehat{\gamma}_0(z)\widehat{\zeta}(z), \\ \widetilde{\gamma}(z) &= \widehat{\gamma}(z)\widehat{\zeta}(z), & \widetilde{\Gamma}_1(z) &= \widehat{\Gamma}_1(z)\widehat{\zeta}(z)\end{aligned}$$

for $z \in \widehat{\Omega} = \Phi(\widetilde{\Omega})$, $\overline{\widetilde{\Omega}} \cap S \neq \emptyset$.

Using the above notation we can rewrite problem (3.28) in the following form in an interior subdomain :

$$\begin{aligned}\eta\tilde{u}_{it} - \nabla_{u_j}T_u^{ij}(\tilde{u}, \tilde{p}_\sigma) &= \eta\tilde{g}_i - \nabla_{u_j}B_u^{ij}(u, \zeta) - T_u^{ij}(u, p_\sigma)\nabla_{u_j}\zeta \\ &\equiv \eta\tilde{g}_i + k_1, \quad i = 1, 2, 3, \\ \widetilde{\eta}_{\sigma t} + \eta\nabla_u \cdot \tilde{u} &= \eta u \cdot \nabla_u \zeta \equiv k_2, \\ \eta c_v(\eta, \Gamma)\widetilde{\gamma}_t - \kappa\nabla_u^2\widetilde{\gamma} + \Gamma p_\Gamma(\eta, \Gamma)\nabla_u \cdot \tilde{u} &= \eta\tilde{k} + \left[\frac{\mu}{2} \sum_{i,j=1}^3 (\xi_{kx_i}\partial_{\xi_k}u_j + \xi_{kx_j}\partial_{\xi_k}u_i)^2 \right. \\ &\quad \left. + (\nu - \mu)(\nabla_u \cdot u)^2 \right] \zeta + \Gamma p_\Gamma(\eta, \Gamma)u \cdot \nabla_u \zeta \\ &\quad - \kappa(\nabla_u^2\zeta\gamma + 2\nabla_u\zeta \cdot \nabla_u\gamma) - \eta c_v(\eta, \Gamma)\zeta\partial_t\theta_{\Omega_t} \\ &\equiv \eta\tilde{k} + k_3,\end{aligned}$$

where $\tilde{p}_\sigma = p_\sigma\zeta$ and

$$\mathbb{B}_u(u, \zeta) = \{B_u^{ij}(u, \zeta)\} = \{\mu(u_i\nabla_{u_j}\zeta + u_j\nabla_{u_i}\zeta) + (\nu - \mu)\delta_{ij}u \cdot \nabla_u\zeta\}.$$

In boundary subdomains we have

$$\begin{aligned}(3.30) \quad \eta\tilde{u}_{it} - \widehat{\nabla}_j\widehat{T}^{ij}(\tilde{u}, \tilde{p}_\sigma) &= \eta\tilde{g}_i - \widehat{\nabla}_j\widehat{B}^{ij}(\widehat{u}, \widehat{\zeta}) - \widehat{T}^{ij}(\widehat{u}, p_\sigma)\widehat{\nabla}_j\widehat{\zeta} \\ &\equiv \eta\tilde{g}_i + k_4^i, \\ \widetilde{\eta}_{\sigma t} + \widehat{\eta}\widehat{\nabla} \cdot \tilde{u} &= \widehat{\eta}\widehat{u} \cdot \widehat{\nabla}\widehat{\zeta} \equiv k_5, \\ \widehat{\eta}c_v(\widehat{\eta}, \widehat{\Gamma})\widetilde{\gamma}_t - \kappa\widehat{\nabla}^2\widetilde{\gamma} + \widehat{\Gamma}p_{\widehat{\Gamma}}(\widehat{\eta}, \widehat{\Gamma})\widehat{\nabla} \cdot \tilde{u} &= \widehat{\eta}\widehat{k} + \left[\frac{\mu}{2} \sum_{i,j=1}^3 (\widehat{\nabla}_i\widehat{u}_j + \widehat{\nabla}_j\widehat{u}_i)^2 + (\nu - \mu)(\widehat{\nabla} \cdot \widehat{u})^2 \right] \widehat{\zeta} \\ &\quad + \widehat{\Gamma}p_{\widehat{\Gamma}}(\widehat{\eta}, \widehat{\Gamma})\widehat{u} \cdot \widehat{\nabla}\widehat{\zeta} - \kappa(\widehat{\nabla}^2\widehat{\zeta}\widehat{\gamma} + \widehat{\nabla}\widehat{\zeta} \cdot \widehat{\nabla}\widehat{\gamma}) \\ &\quad - \widehat{\eta}c_v(\widehat{\eta}, \widehat{\Gamma})\partial_t\theta_{\Omega_t}\widehat{\zeta} \equiv \widehat{\eta}\widehat{k} + k_6, \\ \widehat{\mathbb{T}}(\tilde{u}, \tilde{p}_\sigma)\widehat{n} - \sigma\widehat{\Delta}_{\widehat{S}}\widehat{\xi} \cdot \widehat{n}\widehat{n}\widehat{\zeta} - \sigma\widehat{\Delta}_{\widehat{S}} \int_0^t \tilde{u} dt' \cdot \widehat{n}\widehat{n} &= \frac{2\sigma}{R_0}\widehat{\zeta}\widehat{n} + k_7 + k_8, \\ \widehat{n} \cdot \widehat{\nabla}\widetilde{\gamma} &= \widetilde{\Gamma}_1 + k_9,\end{aligned}$$

where $k_7^i = \widehat{B}^{ij}(\widehat{u}, \widehat{\zeta})\widehat{n}_j$, $k_8 = -\sigma(2\widehat{\nabla} \int_0^t \widehat{u} dt' \widehat{\nabla} \widehat{\zeta} + \int_0^t \widehat{u} dt' \widehat{\nabla}^2 \widehat{\zeta}) \cdot \widehat{n}\widehat{n}$, $k_9 = \widehat{n} \cdot \widehat{\nabla} \widehat{\zeta} \widehat{\gamma}$ and $\widehat{\mathbb{T}}, \widehat{\mathbb{B}}$ indicate that the operator ∇_u is replaced by $\widehat{\nabla}$.

In the considerations below we denote z_1, z_2 by τ and z_3 by n .

LEMMA 3.4. *Let v, ϱ, ϑ_0 be a sufficiently smooth solution of (3.1). Then*

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_x^2 + \frac{p_{\sigma\varrho}}{\varrho} \varrho_x^2 + \frac{\varrho c_v}{\theta} \vartheta_{0x}^2 \right) dx \\ & + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \frac{1}{2} \widetilde{\delta}^{\alpha\beta} \overline{n} \cdot \int_0^t v_{pp^\alpha} dt' \overline{n} \cdot \int_0^t v_{pp^\beta} dt' ds \\ & + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \left| \overline{n} \cdot \int_0^t v_{p^1 p^2} dt' \right|^2 ds \\ & + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \sum_{i=1}^2 \left(\frac{1}{2} \overline{n} \cdot \int_0^t v_{p^i p^i} dt' + 2(H(\cdot, 0) + 2/R_e) \right)^2 ds \\ & + c_0 (\|v_x\|_{1,\Omega_t}^2 + \|\overline{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{0,\Omega_t}^2 + \|\varrho_{\sigma t}\|_{0,\Omega_t}^2 + \|\vartheta_{0xx}\|_{0,\Omega_t}^2) \\ & \leq \varepsilon \left(\|v_{xt}\|_{0,\Omega_t}^2 + \|\vartheta_{0xt}\|_{0,\Omega_t}^2 + \left\| \int_0^t v dt' \right\|_{0,\Omega_t}^2 \right. \\ & \quad \left. + \|H(\cdot, 0) + 2/R_e\|_{0,S^1}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{2,S^1}^2 \right) \\ & + C_1 (\|v\|_{1,0,\Omega_t}^2 + \|\overline{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 + \|\vartheta_{0x}\|_{0,\Omega_t}^2 \\ & \quad + \|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|f\|_{1,\Omega_t}^2 + \|r\|_{1,\Omega_t}^2 + \|\theta_1\|_{2,\Omega_t}^2) \\ & \quad + C_2 (X_3 Y_3 + \|H(\cdot, 0) + 2/R_e\|_{0,S^1}^4), \end{aligned}$$

where the summation over the repeated indices ($\alpha, \beta = 1, 2$) and coordinates ($x, p = (p^1, p^2)$) is assumed, $\widetilde{\delta}^{\alpha\beta}$ on each boundary part $\Sigma_t = S_t \cap \{\zeta(x) \neq 0\}$ (ζ belongs to a partition of unity of Ω_t) is of the form $\widetilde{\delta}^{\alpha\beta} = \delta^{\alpha\beta} + 2\varepsilon^{\alpha\beta}$, $\varepsilon^{\alpha\beta} = -\overline{F}_{p^\alpha} \overline{F}_{p^\beta} (1 + \overline{F}_{p^1}^2 + \overline{F}_{p^2}^2)^{-1}$, \overline{F} is the function such that in the local coordinates $\{y\}$, \sum_t is described by the formula

$$(3.31) \quad y_i = p^i \quad (i = 1, 2), \quad y_3 = \overline{F}(p^1, p^2, t)$$

and $\text{supp } \zeta$ is so small that $|\overline{F}_p| \leq 1/2$. Moreover,

$$\begin{aligned} X_3 &= \|v\|_{2,1,\Omega_t}^2 + |\varrho_\sigma|_{2,1,\Omega_t}^2 + \|\vartheta_0\|_{2,1,\Omega_t}^2 + \|\overline{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2, \\ Y_3 &= X_3 + \|v\|_{3,\Omega_t}^2 + \|\vartheta_{0x}\|_{2,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 + \|\overline{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{3,\Omega_t'}^2 dt'. \end{aligned}$$

P r o o f. Similarly to [16] (see the proof of Lemma 3.4) we obtain the following estimate for interior subdomains:

$$\begin{aligned}
(3.32) \quad & \frac{1}{2} \frac{d}{dt} \int_{\tilde{\Omega}} \left(\eta \tilde{u}_\xi^2 + \frac{p_{\sigma\eta}}{\eta} \tilde{\eta}_{\Omega_t\xi}^2 + \frac{\eta c_v}{\Gamma} \tilde{\gamma}_\xi^2 \right) A d\xi \\
& + \frac{\mu}{2} \|\tilde{u}_\xi\|_{1,\tilde{\Omega}}^2 + \frac{\kappa}{\theta^*} \|\tilde{\gamma}_{\xi\xi}\|_{0,\tilde{\Omega}}^2 + \|\tilde{\eta}_{\Omega_t}\|_{1,\tilde{\Omega}}^2 \\
& \leq \varepsilon (\|\tilde{u}_{\xi\xi}\|_{0,\tilde{\Omega}}^2 + \|\eta_{\sigma\xi}\|_{0,\tilde{\Omega}}^2 + \|\tilde{\gamma}_{\xi\xi}\|_{0,\tilde{\Omega}}^2) \\
& + C_1 (|u|_{1,0,\tilde{\Omega}}^2 + \|v\|_{1,\Omega_t}^2 + \|\gamma_{0\xi}\|_{0,\tilde{\Omega}}^2 + \|\gamma\|_{0,\tilde{\Omega}}^2 \\
& + \|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|\bar{\eta}_{\Omega_t}\|_{0,\tilde{\Omega}}^2 + \|\tilde{g}\|_{0,\tilde{\Omega}}^2 + \|\tilde{k}\|_{0,\tilde{\Omega}}^2) \\
& + C_2 \left[\left(X_3(\tilde{\Omega}) + \int_0^t \|u\|_{3,\tilde{\Omega}}^2 dt' \right) Y_3(\tilde{\Omega}) + \|\gamma\|_{2,\tilde{\Omega}}^2 (\|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|v\|_{1,\Omega_t}^2) \right],
\end{aligned}$$

where

$$\begin{aligned}
X_3(\tilde{\Omega}) &= |u|_{2,1,\tilde{\Omega}}^2 + |\varrho_\sigma|_{2,1,\tilde{\Omega}}^2 + |\gamma_0|_{2,1,\tilde{\Omega}}^2 + \|\bar{\eta}_{\Omega_t}\|_{0,\tilde{\Omega}}^2, \\
Y_3(\tilde{\Omega}) &= X_3(\tilde{\Omega}) + \|u\|_{3,\tilde{\Omega}}^2 + \|\gamma\|_{3,\tilde{\Omega}}^2 + \|\bar{\eta}_{\Omega_t}\|_{0,\tilde{\Omega}}^2 + \int_0^t \|u\|_{3,\tilde{\Omega}}^2 dt'.
\end{aligned}$$

Now, we consider subdomains near the boundary. Differentiate (3.30)₁ with respect to τ , multiply the result by $\tilde{u}_\tau J$ and integrate over $\hat{\Omega}$ (J is the Jacobian of the transformation $x = x(z)$). Next, divide (3.30)₃ by $\hat{\Gamma}$, differentiate the result with respect to τ , multiply by $\tilde{\gamma}_\tau J$ and integrate over $\hat{\Omega}$. Hence using Lemma 5.1 of [18] we get

$$\begin{aligned}
& \frac{1}{2} \frac{d}{dt} \int_{\hat{\Omega}} \left(\hat{\eta} \tilde{u}_\tau^2 + \frac{p_{\sigma\hat{n}}}{\hat{\eta}} \hat{\eta}_{\Omega_t\tau}^2 + \frac{\hat{\eta} c_v}{\hat{\Gamma}} \tilde{\gamma}_\tau^2 \right) J dz + \frac{\mu}{2} \|\tilde{u}_\tau\|_{1,\hat{\Omega}}^2 \\
& + \frac{\kappa}{\theta^*} \|\tilde{\gamma}_{\tau z}\|_{0,\hat{\Omega}}^2 - \int_{\hat{S}} (\hat{n} \hat{\mathbb{T}}(\tilde{u}, \tilde{p}_\sigma))_{,\tau} \tilde{u}_\tau J dz' - \kappa \int_{\hat{S}} \left(\hat{n} \frac{1}{\hat{\Gamma}} \hat{\nabla} \tilde{\gamma} \right)_{,\tau} \tilde{\gamma}_\tau J dz' \\
& \leq \varepsilon (\|\tilde{u}_{zz}\|_{0,\hat{\Omega}}^2 + \|\hat{\eta}_{\sigma z}\|_{0,\hat{\Omega}}^2 + \|\hat{\gamma}_{0zz}\|_{0,\hat{\Omega}}^2) \\
& + C_1 (|\hat{u}|_{1,0,\hat{\Omega}}^2 + \|v\|_{1,\Omega_t}^2 + \|\hat{\gamma}_{0\tau}\|_{0,\hat{\Omega}}^2 \\
& + \|\hat{\gamma}\|_{0,\hat{\Omega}}^2 + \|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|\bar{\eta}_{\Omega_t}\|_{0,\hat{\Omega}}^2 + \|\tilde{g}\|_{1,\hat{\Omega}}^2 + \|\tilde{k}\|_{1,\hat{\Omega}}^2) \\
& + C_2 \left[\left(X_2(\hat{\Omega}) + \int_0^t \|\hat{u}\|_{3,\hat{\Omega}}^2 dt' \right) Y_2(\hat{\Omega}) + \|\hat{\gamma}\|_{2,\hat{\Omega}}^2 (\|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|v\|_{1,\Omega_t}^2) \right],
\end{aligned}$$

where $X_2(\widehat{\Omega})$ and $Y_2(\widehat{\Omega})$ are defined analogously to $X_2(\widetilde{\Omega})$ and $Y_2(\widetilde{\Omega})$.

Using the boundary condition (3.30)₄ we have

$$\begin{aligned}
 (3.33) \quad & - \int_{\widehat{S}} (\widehat{n} \widehat{\mathbb{T}}(\widetilde{u}, \widetilde{p}_\sigma)),_\tau \widetilde{u}_\tau J dz' \\
 & \leq - \frac{\sigma}{2} \frac{d}{dt} \int_{\widehat{S}} g^{\alpha\beta} \widehat{n} \cdot \int_0^t \widetilde{u}_{pp^\alpha} dt' \widehat{n} \cdot \int_0^t \widetilde{u}_{pp^\beta} dt' J dz' \\
 & \quad - \sigma \int_{\widehat{S}} (\widehat{H}(\cdot, 0) + 2/R_e) \widehat{\zeta} \cdot \widetilde{u}_{pp} \cdot \widehat{n} J dz' \\
 & \quad + \varepsilon \left(\left\| \int_0^t \widetilde{u} dt' \right\|_{2, \widehat{S}}^2 + \|\widetilde{u}_{zz}\|_{0, \widehat{\Omega}}^2 + \|(\widehat{H}(\cdot, 0) + 2/R_e) \widehat{\zeta}\|_{0, \widehat{S}}^2 \right. \\
 & \quad \left. + \|R(\cdot, t) - R(\cdot, 0)\|_{2, S^1}^2 \right) \\
 & \quad + C_2 \left(\|\widehat{u}\|_{0, \widehat{\Omega}}^2 + \|\widehat{u}\|_{2, \widehat{\Omega}}^2 \left\| \int_0^t \widetilde{u} dt' \right\|_{3, \widehat{\Omega}}^2 \right).
 \end{aligned}$$

By the boundary condition (3.30)₅ we get

$$\begin{aligned}
 (3.34) \quad & - \kappa \int_{\widehat{S}} \left(\widehat{n} \cdot \frac{1}{\widehat{\Gamma}} \widehat{\nabla} \widehat{\gamma} \right),_\tau \widetilde{\gamma}_\tau J dz' \\
 & \leq \varepsilon \|\widehat{\gamma}_{0zz}\|_{0, \widehat{\Omega}}^2 + C_1 (\|\widehat{\gamma}\|_{0, \widehat{\Omega}}^2 + \|\widehat{\gamma}_{0z}\|_{0, \widehat{\Omega}}^2 + \|\widetilde{\Gamma}_1\|_{2, \widehat{\Omega}}^2) \\
 & \quad + C_2 \|\widehat{\gamma}\|_{2, \widehat{\Omega}}^2 \left(\|\widehat{\gamma}_0\|_{2, \widehat{\Omega}}^2 + \|\widehat{\gamma}\|_{2, \widehat{\Omega}}^2 + \|\widehat{\eta}_\sigma\|_{2, \widehat{\Omega}}^2 + \left\| \int_0^t \widehat{u} dt' \right\|_{3, \widehat{\Omega}}^2 \right).
 \end{aligned}$$

To obtain (3.33) and (3.34) we have applied the interpolation inequality (2.2) (see Lemma 2.1).

For the quantities

$$\begin{aligned}
 & \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \frac{p_\sigma \widehat{\eta}}{\widehat{\eta}} \widetilde{\eta}_{\Omega_t n} J dz + c_0 \|\widetilde{\eta}_{\Omega_t n}\|_{0, \widehat{\Omega}}^2, \\
 & \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \widehat{\eta} \widetilde{u}_{3n}^2 J dz + c_0 \|\widetilde{u}_{3nn}\|_{0, \widehat{\Omega}}^2, \\
 & \|\widetilde{\eta}_{\Omega_t}\|_{0, \Omega_t}^2, \quad \|\widetilde{u}'_{z\tau}\|_{0, \widehat{\Omega}}^2, \quad \|\widetilde{\eta}_{\Omega_t \tau}\|_{0, \widehat{\Omega}}^2, \quad \|\widetilde{u}'_{nn}\|_{0, \widehat{\Omega}}^2, \quad \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \widehat{\eta} \widetilde{u}_n^2 J dz, \\
 & \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \frac{\widehat{\eta} c_v}{\widehat{\Gamma}} \widetilde{\gamma}_n^2 J dz + \frac{\kappa}{\theta^*} \|\widetilde{\gamma}_{nn}\|_{0, \widehat{\Omega}}^2
 \end{aligned}$$

we obtain the same estimates as in the proof of Lemma 3.4 of [16]. Therefore, we have

$$\begin{aligned}
(3.35) \quad & \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \left(\widehat{\eta} \widetilde{u}_z^2 + \frac{p_{\sigma\eta}}{\widehat{\eta}} \widetilde{\eta}_{\Omega_t z}^2 + \frac{\widehat{\eta} c_v}{\widehat{\Gamma}} \widetilde{\gamma}_z^2 \right) J dz \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{\widehat{S}} \left[g^{\alpha\beta} \widehat{n} \cdot \int_0^t \widetilde{u}_{pp^\alpha} dt' \widehat{n} \cdot \int_0^t \widetilde{u}_{pp^\beta} dt' \right. \\
& \quad \left. + 2(\widehat{H}(\cdot, 0) + 2/R_e) \widehat{\zeta} \widehat{n} \cdot \int_0^t \widetilde{u}_{pp} dt' \right] J dz' + \frac{\mu}{2} \|\widetilde{u}_z\|_{1,\widehat{\Omega}}^2 \\
& + \frac{\kappa}{\theta_*} \|\widetilde{\gamma}_{zz}\|_{0,\widehat{\Omega}}^2 + c_0 \|\widetilde{\eta}_{\Omega_t}\|_{1,\widehat{\Omega}}^2 \\
& \leq \varepsilon \left(\|\widetilde{u}_{zz}\|_{0,\widehat{\Omega}}^2 + \|\widehat{\eta}_{\sigma z}\|_{0,\widehat{\Omega}}^2 + \|\widehat{\gamma}_{0zz}\|_{0,\widehat{\Omega}}^2 + \|\widetilde{u}_{zt}\|_{0,\widehat{\Omega}}^2 + \|\widetilde{\gamma}_{0zt}\|_{0,\widehat{\Omega}}^2 \right. \\
& \quad \left. + \left\| \int_0^t \widetilde{u} dt' \right\|_{2,\widehat{S}}^2 + \|(\widehat{H}(\cdot, 0) + 2/R_e) \widehat{\zeta}\|_{0,\widehat{S}}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{2,S^1}^2 \right) \\
& + C_1 (\|\widehat{u}\|_{1,0,\widehat{\Omega}}^2 + \|v\|_{1,\Omega_t}^2 + \|\widehat{\gamma}_{0\tau}\|_{0,\widehat{\Omega}}^2 + \|\widehat{\gamma}\|_{0,\widehat{\Omega}}^2 \\
& + \|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|\widehat{\eta}_{\Omega_t}\|_{0,\widehat{\Omega}}^2 + \|\widetilde{g}\|_{1,\widehat{\Omega}}^2 + \|\widetilde{k}\|_{1,\widehat{\Omega}}^2) \\
& + C_2 \left[\left(X_2(\widehat{\Omega}) + \int_0^t \|\widetilde{u}\|_{3,\widehat{\Omega}}^2 dt' \right) Y_2(\widehat{\Omega}) + \|\widehat{\gamma}\|_{2,\widehat{\Omega}}^2 (\|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|v\|_{1,\Omega_t}^2) \right].
\end{aligned}$$

We estimate the second term on the left-hand side of (3.35) in the same way as in the proof of Lemma 4.4 of [19]. Going back to the variables ξ in (3.35), next from the resulting estimate and (3.32), after summing over all neighbourhoods of the partition of unity and finally going back to the variables x and using (3.26) we get

$$\begin{aligned}
(3.36) \quad & \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_x^2 + \frac{p_{\sigma\varrho}}{\varrho} \varrho_{\sigma x}^2 + \frac{\varrho c_v}{\theta} \vartheta_{0x}^2 \right) dx \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \frac{1}{2} \widetilde{\delta}^{\alpha\beta} \overline{n} \cdot \int_0^t v_{pp^\alpha} dt' \overline{n} \cdot \int_0^t v_{pp^\beta} dt' ds \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \left| \overline{n} \cdot \int_0^t v_{p^1 p^2} dt' \right|^2 ds \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \sum_{i=1}^2 \left(\frac{1}{2} \overline{n} \cdot \int_0^t v_{p^i p^i} dt' + 2(\widehat{H}(\cdot, 0) + 2/R_e) \right)^2 ds
\end{aligned}$$

$$\begin{aligned}
& + c_0 (\|v_x\|_{1,\Omega_t}^2 + \|\vartheta_{0xx}\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{0,\Omega_t}^2 + \|\varrho_{\sigma t}\|_{0,\Omega_t}^2) \\
& \leq \varepsilon \left(\|v_{xt}\|_{0,\Omega_t}^2 + \|\vartheta_{0xt}\|_{0,\Omega_t}^2 + \left\| \int_0^t v dt' \right\|_{2,S_t}^2 \right. \\
& \quad \left. + \|H(\cdot, 0) + 2/R_e\|_{0,S_t}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{2,S^1}^2 \right) \\
& \quad + C_1 (|v|_{1,0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 + \|\vartheta_{0x}\|_{0,\Omega_t}^2 \\
& \quad + \|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|f\|_{1,\Omega_t}^2 + \|r\|_{1,\Omega_t}^2 + \|\theta_1\|_{2,\Omega_t}^2) \\
& \quad + C_2 X_3 Y_3 + 4\sigma \frac{d}{dt} \int_{S_t} (H(\cdot, 0) + 2/R_e)^2 ds.
\end{aligned}$$

In virtue of the interpolation inequality (2.2) we have

$$\begin{aligned}
(3.37) \quad & \left| \frac{d}{dt} \int_{S_t} (H(\cdot, 0) + 2/R_e)^2 ds \right| \\
& \leq \varepsilon \|v_{xx}\|_{0,\Omega_t}^2 + C_1 \|v\|_{0,\Omega_t}^2 + C_2 \|H(\cdot, 0) + 2/R_e\|_{0,S^1}^4.
\end{aligned}$$

Writing the boundary condition (3.1)₄ locally we obtain

$$(3.38) \quad \sigma \widehat{\Delta}_{\widehat{S}} \int_0^t \widetilde{u} dt' = -\sigma \left(\widehat{\Delta}_{\widehat{S}} \widehat{\xi} + \frac{2}{R_e} \widehat{n} \right) \widehat{\zeta} - \widehat{\mathbb{T}}_u(\widetilde{u}, \widetilde{p}_\sigma) \widehat{n} + I_1 + I_2,$$

where

$$I_1^i = -\widehat{B}^{ij}(\widehat{u}, \widehat{\zeta}) \widehat{n}_j, \quad I_2 = \sigma \left(2 \widehat{\nabla} \int_0^t \widehat{u} dt' \widehat{\nabla} \widehat{\zeta} + \int_0^t \widehat{u} dt' \widehat{\nabla}^2 \widehat{\zeta} \right).$$

Multiply (3.38) by $\int_0^t \widetilde{u} dt'$, next differentiate with respect to τ and multiply by $\int_0^t \widetilde{u}_\tau dt'$. Integrating the result over \widehat{S} and summing over all neighbourhoods of the partition of unity we get

$$\begin{aligned}
(3.39) \quad & \left\| \int_0^t v dt' \right\|_{2,S_t}^2 \\
& \leq \varepsilon (\|v\|_{2,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{0,\Omega_t}^2 + \|\vartheta_{0x}\|_{0,\Omega_t}^2) \\
& \quad + C_1 (\|v\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 \\
& \quad + \left\| \int_0^t v dt' \right\|_{0,\Omega_t}^2 + \|H(\cdot, 0) + 2/R_e\|_{0,S^1}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{2,S^1}^2) \\
& \quad + C_2 (\|v\|_{2,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{1,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\vartheta_0\|_{2,\Omega_t}^2) \int_0^t \|v\|_{3,\Omega_{t'}}^2 dt'.
\end{aligned}$$

From (3.36), (3.37) and (3.39) we obtain (3.31). ■

Now, we formulate Lemmas 3.5–3.7, the proofs of which are similar to the proofs of Lemmas 3.5–3.7 of [16]. The boundary terms associated with the boundary condition (3.1)₄ are estimated in the same way as in Lemmas 4.5–4.7 of [19] and similarly to Lemmas 4.1, 4.2 and 4.4.

LEMMA 3.5. *Let v, ϱ, ϑ_0 be a sufficiently smooth solution of problem (3.1). Then*

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_{xx}^2 + \frac{p_{\sigma\varrho}}{\varrho} \varrho_{\sigma xx}^2 + \frac{\varrho c_v}{\theta} \vartheta_{0xx}^2 \right) dx \\ & + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \frac{1}{2} \tilde{\delta}^{\alpha\beta} \bar{n} \cdot \int_0^t v_{p^\gamma p^\delta p^\alpha} dt' \bar{n} \cdot \int_0^t v_{p^\gamma p^\delta p^\beta} dt' ds \\ & + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \left| \bar{n} \cdot \int_0^t v_{p^1 p^2 p} dt' \right|^2 ds \\ & + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \sum_{i=1}^2 \left(\frac{1}{2} \bar{n} \cdot \int_0^t v_{pp^i p^i} dt' + 2(H(\cdot, 0) + 2/R_e)_{,p} \right)^2 ds \\ & + c_0 (\|v_{xx}\|_{1,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{1,\Omega_t}^2 + \|\vartheta_{0xxx}\|_{0,\Omega_t}^2) \\ & \leq \varepsilon (\|v_{xxt}\|_{0,\Omega_t}^2 + \|\vartheta_{0xxt}\|_{0,\Omega_t}^2 + \|H(\cdot, 0) + 2/R_e\|_{1,S^1}^2 \\ & + \|R(\cdot, t) - R(\cdot, 0)\|_{3,S^1}^2) \\ & + C_1 (\|v\|_{2,1,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{0,\Omega_t}^2 + \|\varrho_{\sigma t}\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\vartheta_{0x}\|_{1,\Omega_t}^2 \\ & + \|\vartheta_{0t}\|_{1,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 + \|f\|_{1,\Omega_t}^2 + \|r\|_{1,\Omega_t}^2 + \|\theta_1\|_{3,\Omega_t}^2) \\ & + C_2 [X_4 (1 + X_4) Y_4 + \|H(\cdot, 0) + 2/R_e\|_{1,S^1}^4], \end{aligned}$$

where the summation over the repeated indices $(\alpha, \beta, \gamma, \delta)$ and coordinates $x, p = (p^1, p^2)$ is assumed and

$$\begin{aligned} X_4 &= \|v\|_{3,\Omega_t}^2 + \|v_t\|_{1,\Omega_t}^2 + |\varrho_\sigma|_{2,1,\Omega_t}^2 + |\vartheta_0|_{2,1,\Omega_t}^2 + \|\vartheta_0\|_{3,\Omega_t}^2 \\ &+ \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{3,\Omega_{t'}}^2 dt', \\ Y_4 &= \|v\|_{4,\Omega_t}^2 + \|v_t\|_{1,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{2,\Omega_t}^2 + |\varrho_\sigma|_{2,1,\Omega_t}^2 + \|\vartheta_{0x}\|_{3,\Omega_t}^2 \\ &+ \|\vartheta_{0t}\|_{1,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{4,\Omega_{t'}}^2 dt'. \blacksquare \end{aligned}$$

LEMMA 3.6. *Let v, ϱ, ϑ_0 be a sufficiently smooth solution of problem (3.1). Then*

$$\frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_{xt}^2 + \frac{p_{\sigma\varrho}}{\varrho} \varrho_{\sigma xt}^2 + \frac{\varrho c_v}{\theta} \vartheta_{0xt}^2 \right) dx + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} g^{\alpha\beta} \bar{n} \cdot v_{pp^\alpha} \bar{n} \cdot v_{pp^\beta} ds$$

$$\begin{aligned}
& + c_0 (\|v_t\|_{2,\Omega_t}^2 + \|\varrho_{\sigma t}\|_{1,\Omega_t}^2 + \|\vartheta_{0xxt}\|_{0,\Omega_t}^2) \\
& \leq \varepsilon (\|v_{xtt}\|_{0,\Omega_t}^2 + \|\vartheta_{0xtt}\|_{0,\Omega_t}^2 + \|H(\cdot, 0) + 2/R_e\|_{0,S^1}^2) \\
& \quad + C_1 (\|v\|_{2,0,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{0,\Omega_t}^2 + \|\varrho_{\sigma t}\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 \\
& \quad + \|\vartheta_{0x}\|_{1,\Omega_t}^2 + \|\vartheta_{0t}\|_{1,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 + |f|_{1,0,\Omega_t}^2 + |r|_{1,0,\Omega_t}^2 \\
& \quad + \|\theta_{1t}\|_{2,\Omega_t}^2 + \|\theta_1\|_{1,\Omega_t}^2) + C_2 X_5 (1 + X_5) Y_5,
\end{aligned}$$

where to describe S_t we have used formula (3.31) and

$$\begin{aligned}
X_5 &= |v|_{3,1,\Omega_t}^2 + |\varrho_\sigma|_{2,0,\Omega_t}^2 + |\vartheta_0|_{3,1,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{3,\Omega_{t'}}^2 dt', \\
Y_5 &= |v|_{4,2,\Omega_t}^2 + |\varrho_\sigma|_{3,1,\Omega_t}^2 + |\vartheta_{0t}|_{3,2,\Omega_t}^2 + \|\vartheta_{0x}\|_{3,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 \\
&\quad + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{4,\Omega_{t'}}^2 dt'.
\end{aligned}$$

LEMMA 3.7. Let v, ϱ, ϑ_0 be a sufficiently smooth solution of problem (3.1).

Then

$$\begin{aligned}
& \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_{tt}^2 + \frac{p_{\sigma\varrho}}{\varrho} \varrho_{\sigma tt}^2 + \frac{\varrho c_v}{\theta} \vartheta_{0tt}^2 \right) dx + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} g^{\alpha\beta} \bar{n} \cdot v_{s^\alpha t} \bar{n} \cdot v_{s^\beta t} ds \\
& \quad + c_0 (\|v_{tt}\|_{1,\Omega_t}^2 + \|\varrho_{\sigma tt}\|_{0,\Omega_t}^2 + \|\vartheta_{0tt}\|_{1,\Omega_t}^2) \\
& \leq \varepsilon \|H(\cdot, 0) + 2/R_e\|_{0,S^1}^2 + C_1 (\|v_t\|_{1,\Omega_t}^2 + \|\vartheta_{0t}\|_{1,\Omega_t}^2 + |f|_{1,0,\Omega_t}^2 + \|f_{tt}\|_{0,\Omega_t}^2 \\
& \quad + |r|_{1,0,\Omega_t}^2 + \|r_{tt}\|_{0,\Omega_t}^2 + \|r\|_{0,\Omega_t}^2 + \|\theta_1\|_{3,1,\Omega_t}^2) + C_2 X_6 (1 + X_6) Y_6,
\end{aligned}$$

where

$$\begin{aligned}
X_6 &= |v|_{3,1,\Omega_t}^2 + |\varrho_\sigma|_{2,0,\Omega_t}^2 + |\vartheta_0|_{3,1,\Omega_t}^2, \\
Y_6 &= |v|_{4,2,\Omega_t}^2 + |\varrho_\sigma|_{3,1,\Omega_t}^2 + |\vartheta_0|_{4,2,\Omega_t}^2. \blacksquare
\end{aligned}$$

Summarizing, from Lemmas 3.5–3.7 we obtain

LEMMA 3.8. Let v, ϱ, ϑ_0 be a sufficiently smooth solution of problem (3.1).

Then

$$\begin{aligned}
& \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho |D_{x,t}^2 v|^2 + \frac{p_{\sigma\varrho}}{\varrho} |D_{x,t}^2 \varrho_\sigma|^2 + \frac{\varrho c_v}{\theta} |D_{x,t}^2 \vartheta_0|^2 \right) dx \\
& \quad + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \frac{1}{2} \tilde{\delta}^{\alpha\beta} \bar{n} \cdot \int_0^t v_{p^\gamma p^\delta p^\alpha} dt' \bar{n} \cdot \int_0^t v_{p^\gamma p^\delta p^\beta} dt' ds \\
& \quad + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \left| \bar{n} \cdot \int_0^t v_{pp^1 p^2} dt' \right|^2 ds
\end{aligned}$$

$$\begin{aligned}
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \sum_{i=1}^2 \left(\frac{1}{2} \bar{n} \cdot \int_0^t v_{p^1 p^2 p} dt' + 2(H(\cdot, 0) + 2/R_e)_{,p} \right)^2 ds \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} g^{\alpha\beta} \bar{n} \cdot v_{pp^\alpha} \bar{n} \cdot v_{pp^\beta} ds + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} g^{\alpha\beta} \bar{n} \cdot v_{s^\alpha t} \bar{n} \cdot v_{s^\beta t} ds \\
& + c_0 (|v|_{3,1,\Omega_t}^2 + |\varrho_{\sigma t}|_{1,0,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{0,\Omega_t}^2 + \|\vartheta_{0t}\|_{2,1,\Omega_t}^2 + \|\vartheta_{0xxx}\|_{0,\Omega_t}^2) \\
& \leq \varepsilon (\|H(\cdot, 0) + 2/R_e\|_{1,S^1}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{3,S^1}^2) \\
& + C_1 (|v|_{2,0,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{0,\Omega_t}^2 + \|\varrho_{\sigma t}\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 \\
& + \|\vartheta_{0x}\|_{1,\Omega_t}^2 + \|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 + |f|_{1,0,\Omega_t}^2 + \|f_{tt}\|_{0,\Omega_t}^2 \\
& + |r|_{1,0,\Omega_t}^2 + \|r_{tt}\|_{0,\Omega_t}^2 + |\theta_1|_{3,1,\Omega_t}^2) \\
& + C_2 [X_7 (1 + X_7) Y_7 + \|H(\cdot, 0) + 2/R_e\|_{1,S^1}^4],
\end{aligned}$$

where

$$\begin{aligned}
X_7 & = |v|_{3,1,\Omega_t}^2 + |\varrho_\sigma|_{2,0,\Omega_t}^2 + |\vartheta_0|_{3,1,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{3,\Omega_{t'}}^2 dt', \\
Y_7 & = |v|_{4,2,\Omega_t}^2 + |\varrho_\sigma|_{3,1,\Omega_t}^2 + |\vartheta_{0t}|_{3,2,\Omega_t}^2 + \|\vartheta_{0x}\|_{3,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 \\
& + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{4,\Omega_{t'}}^2 dt'. \blacksquare
\end{aligned}$$

LEMMA 3.9. Let v, ϱ, ϑ_0 be a sufficiently smooth solution of problem (3.1).

Then

$$\begin{aligned}
(3.40) \quad & \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_{xxx}^2 + \frac{p_\sigma \varrho}{\varrho} \varrho_{\sigma xxx}^2 + \frac{\varrho c_v}{\theta} \vartheta_{0xxx}^2 \right) dx \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \frac{1}{2} \tilde{\delta}^{\alpha\beta} \bar{n} \cdot \int_0^t v_{p^\gamma p^\delta p^\eta p^\alpha} dt' \bar{n} \cdot \int_0^t v_{p^\gamma p^\delta p^\eta p^\beta} dt' ds \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \left| \bar{n} \cdot \int_0^t v_{ppp^1 p^2} dt' \right|^2 ds \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \sum_{i=1}^2 \left(\frac{1}{2} \bar{n} \cdot \int_0^t v_{ppp^i p^i} dt' + 2(H(\cdot, 0) + 2/R_e)_{,pp} \right)^2 ds \\
& + c_0 (\|v_{xxx}\|_{1,\Omega_t}^2 + \|\varrho_{\sigma xxx}\|_{0,\Omega_t}^2 + \|\vartheta_{0xxxx}\|_{0,\Omega_t}^2) \\
& \leq \varepsilon (\|v_{xxxt}\|_{0,\Omega_t}^2 + \|\vartheta_{0xxxt}\|_{0,\Omega_t}^2 + \|H(\cdot, 0) + 2/R_0\|_{2,S^1}^2 \\
& + \|R(\cdot, t) - R(\cdot, 0)\|_{4,S^1}^2) + C_1 (|v|_{3,2,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{1,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 \\
& + \|\vartheta_{0x}\|_{2,\Omega_t}^2 + \|\vartheta_{0t}\|_{2,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 + \|f\|_{2,\Omega_t}^2 + \|r\|_{2,\Omega_t}^2)
\end{aligned}$$

$$\begin{aligned}
& + \|\theta_1\|_{4,\Omega_t}^2 + \left\| \int_0^t v dt' \right\|_{0,\Omega_t}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{0,S^1}^2 \Big) \\
& + C_2 \left[\|H(\cdot, 0) + 2/R_0\|_{2,S^1}^4 + \|R(\cdot, t) - R(\cdot, 0)\|_{3,S^1}^2 \left\| \int_0^t v dt' \right\|_{4,S_t}^2 \right. \\
& \left. + \|R(\cdot, t) - R(\cdot, 0)\|_{4+1/2,S^1}^2 \left\| \int_0^t v dt' \right\|_{3,S_t}^2 + X_8(1 + X_8^2)Y_8 \right],
\end{aligned}$$

where the summation over repeated indices ($\alpha, \beta, \gamma, \delta = 1, 2$) and coordinates $(x, p = (p^1, p^2), i = 1, 2)$ is assumed and

$$\begin{aligned}
X_8 &= |v|_{3,2,\Omega_t}^2 + |\varrho_\sigma|_{3,2,\Omega_t}^2 + |\vartheta_0|_{3,2,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{3,\Omega_{t'}}^2 dt', \\
Y_8 &= |v|_{4,3,\Omega_t}^2 + |\varrho_\sigma|_{3,2,\Omega_t}^2 + \|\vartheta_{0x}\|_{3,\Omega_t}^2 + \|\vartheta_{0t}\|_{3,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 \\
&\quad + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{4,\Omega_{t'}}^2 dt'.
\end{aligned}$$

P r o o f. For interior subdomains we obtain the estimate (see [16], proof of Lemma 3.9)

$$\begin{aligned}
(3.41) \quad & \frac{1}{2} \frac{d}{dt} \int_{\tilde{\Omega}} \left(\eta \tilde{u}_{\xi\xi\xi\xi}^2 + \frac{p_{\sigma\eta}}{\eta} \tilde{\eta}_{\Omega_t\xi\xi\xi\xi}^2 + \frac{\eta c_v}{\Gamma} \tilde{\gamma}_{\xi\xi\xi\xi}^2 \right) A d\xi + \frac{\mu}{2} \|\tilde{u}_{\xi\xi\xi\xi}\|_{1,\tilde{\Omega}}^2 \\
& + \|\tilde{\eta}_{\Omega_t\xi\xi\xi\xi}\|_{0,\tilde{\Omega}}^2 + \frac{\kappa}{\theta^*} \|\tilde{\gamma}_{\xi\xi\xi\xi}\|_{0,\tilde{\Omega}}^2 \\
& \leq \varepsilon (\|\tilde{u}_{\xi\xi\xi\xi}\|_{0,\tilde{\Omega}}^2 + \|\tilde{\gamma}_{\xi\xi\xi\xi}\|_{0,\tilde{\Omega}}^2 + \|\tilde{\eta}_{\sigma\xi\xi\xi\xi}\|_{0,\tilde{\Omega}}^2) \\
& + C_1 (|u|_{3,2,\tilde{\Omega}}^2 + \|\gamma_{0\xi}\|_{2,\tilde{\Omega}}^2 + \|\gamma\|_{0,\tilde{\Omega}}^2 + \|\eta_{\sigma\xi}\|_{1,\tilde{\Omega}}^2 \\
& + \|\bar{\eta}_{\Omega_t}\|_{0,\tilde{\Omega}}^2 + \|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|v\|_{1,\Omega_t}^2 + \|\tilde{g}\|_{2,\tilde{\Omega}}^2 + \|\tilde{k}\|_{2,\tilde{\Omega}}^2) \\
& + C_2 \left[\left(X_8(\tilde{\Omega}) + \int_0^t \|u\|_{3,\tilde{\Omega}}^2 dt' \right) (1 + X_8^2(\tilde{\Omega})) Y_8(\tilde{\Omega}) \right. \\
& \left. + \|\gamma\|_{4,\tilde{\Omega}}^2 (\|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|v\|_{1,\Omega_t}^2) \right],
\end{aligned}$$

where

$$\begin{aligned}
X_8(\tilde{\Omega}) &= |u|_{3,2,\tilde{\Omega}}^2 + |\bar{\eta}_{\Omega_t}|_{3,2,\tilde{\Omega}}^2 + |\gamma|_{3,2,\tilde{\Omega}}^2 + |\eta_\sigma|_{3,2,\tilde{\Omega}}^2 + |\gamma_0|_{3,2,\tilde{\Omega}}^2, \\
Y_8(\tilde{\Omega}) &= |u|_{4,3,\tilde{\Omega}}^2 + |\bar{\eta}_{\Omega_t}|_{3,2,\tilde{\Omega}}^2 + |\gamma|_{4,3,\tilde{\Omega}}^2 + \|\eta_{\sigma t}\|_{2,\tilde{\Omega}}^2 \\
& + \|\gamma_{0t}\|_{3,\tilde{\Omega}}^2 + \int_0^t \|u\|_{4,\tilde{\Omega}}^2 dt.
\end{aligned}$$

For boundary subdomains we have

$$\begin{aligned}
& \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \left(\widehat{\eta} \widetilde{u}_{\tau\tau\tau}^2 + \frac{p_{\sigma\eta}}{\widehat{\eta}} \widetilde{\eta}_{\Omega_t\tau\tau\tau}^2 + \frac{\widehat{\eta} c_v}{\widehat{I}} \widetilde{\gamma}_{\tau\tau\tau}^2 \right) J dz \\
& + \frac{\mu}{2} \|\widetilde{u}_{\tau\tau\tau}\|_{1,\widehat{\Omega}}^2 + \frac{\kappa}{\theta^*} \|\widetilde{\gamma}_{\tau\tau\tau z}\|_{0,\widehat{\Omega}}^2 \\
& - \int_{\widehat{S}} (\widehat{n} \widehat{\mathbb{T}}(\widetilde{u}, \widetilde{p}_\sigma))_{,\tau\tau\tau} \widetilde{u}_{\tau\tau\tau} J dz' - \kappa \int_{\widehat{S}} \left(\widehat{n} \frac{1}{\widehat{I}} \widehat{\nabla} \widetilde{\gamma} \right)_{,\tau\tau\tau} \widetilde{\gamma}_{\tau\tau\tau} J dz' \\
& \leq \varepsilon (\|\widehat{u}_{zzzz}\|_{0,\widehat{\Omega}}^2 + \|\widehat{\gamma}_{0zzzz}\|_{0,\widehat{\Omega}}^2 + \|\widehat{\eta}_{\sigma zzzz}\|_{0,\widehat{\Omega}}^2) \\
& + C_1 (\|\widehat{u}\|_{3,2,\widehat{\Omega}}^2 + \|\widehat{\gamma}_{0z}\|_{2,\widehat{\Omega}}^2 + \|\widehat{\gamma}\|_{0,\widehat{\Omega}}^2 + \|\widehat{\eta}_{\sigma z}\|_{1,\widehat{\Omega}}^2 + \|\overline{\eta}_{\Omega_t}\|_{0,\widehat{\Omega}}^2) \\
& + \|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|v\|_{1,\Omega_t}^2 + \|\widetilde{g}\|_{2,\widehat{\Omega}}^2 + \|\widetilde{k}\|_{2,\widehat{\Omega}}^2 + \|\widetilde{I}_1\|_{4,\widehat{\Omega}}^2 \\
& + C_2 \left[\left(X_8(\widehat{\Omega}) + \int_0^t \|\widehat{u}\|_{3,\widehat{\Omega}}^2 dt' \right) (1 + X_8^2(\widehat{\Omega})) Y_8(\widehat{\Omega}) \right. \\
& \left. + \|\widehat{\gamma}\|_{4,\widehat{\Omega}}^2 (\|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|v\|_{1,\Omega_t}^2) \right],
\end{aligned}$$

where $X_8(\widehat{\Omega})$ and $Y_8(\widehat{\Omega})$ are defined analogously to $X_8(\widetilde{\Omega})$ and $Y_8(\widetilde{\Omega})$. The boundary conditions (3.30)₄ and (3.30)₅ yield

$$\begin{aligned}
& - \int_{\widehat{S}} (\widehat{n} \widehat{\mathbb{T}}(\widetilde{u}, \widetilde{p}_\sigma))_{,\tau\tau\tau} \widetilde{u}_{\tau\tau\tau} J dz' \\
& = - \sigma \int_{\widehat{S}} (\widehat{\Delta}_{\widehat{S}} \widehat{\xi} \cdot \widehat{n} \widehat{n} \widehat{\zeta} + (2/R_e) \widehat{n} \widehat{\zeta})_{,\tau\tau\tau} \widetilde{u}_{\tau\tau\tau} dz' \\
& - \sigma \int_{\widehat{S}} \left(\widehat{\Delta}_{\widehat{S}} \int_0^t \widetilde{u} d\tau \cdot \widehat{n} \widehat{n} \right)_{,\tau\tau\tau} \widetilde{u}_{\tau\tau\tau} J dz' + \int_{\widehat{S}} (k_7 + k_8)_{,\tau\tau\tau} \widetilde{u}_{\tau\tau\tau} J dz' \\
& \leq - \frac{\sigma}{2} \frac{d}{dt} \int_{\widehat{S}} g^{\alpha\beta} \overline{n} \cdot \int_0^t \widetilde{u}_{\tau\tau\tau p^\alpha} dt' \overline{n} \cdot \int_0^t \widetilde{u}_{\tau\tau\tau p^\beta} dt' J dz' \\
& - \frac{\sigma}{2} \frac{d}{dt} \int_{\widehat{S}} (\widehat{H}(\cdot, 0) + 2/R_e)_{,\tau\tau} \widehat{\zeta} \widehat{n} \cdot \int_0^t \widetilde{u}_{\tau\tau\tau_i \tau_i} dt' J dz' \\
& + \varepsilon \left(\|\widetilde{u}_{\tau\tau\tau}\|_{1,\widehat{\Omega}}^2 + \left\| \int_0^t \widetilde{u} dt' \right\|_{4,\widehat{S}}^2 + \|\widehat{H}(\cdot, 0) + 2/R_e\|_{2,\widehat{S}}^2 \right. \\
& \left. + \|R(\cdot, t) - R(\cdot, 0)\|_{4,S^1}^2 \right) \\
& + C_1 \left(\|\widehat{u}\|_{3,\widehat{\Omega}}^2 + \left\| \int_0^t \widetilde{u} dt' \right\|_{0,\widehat{\Omega}}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{0,S^1}^2 \right)
\end{aligned}$$

$$\begin{aligned}
& + C_2 \left(\|R(\cdot, t) - R(\cdot, 0)\|_{3, S^1}^2 \left\| \int_0^t \widehat{u} dt' \right\|_{4, \widehat{S}}^2 \right. \\
& \quad \left. + \|R(\cdot, t) - R(\cdot, 0)\|_{4+1/2, S^1}^2 \left\| \int_0^t \widehat{u} dt' \right\|_{3, \widehat{S}}^2 + \|\widehat{u}\|_{3, \widehat{\Omega}}^2 \left\| \int_0^t \widehat{u} dt' \right\|_{4, \widehat{S}}^2 \right),
\end{aligned}$$

and

$$\begin{aligned}
& -\kappa \int_{\widehat{S}} \left(\widehat{n} \cdot \frac{1}{\widehat{\Gamma}} \widehat{\nabla} \widehat{\gamma} \right)_{,\tau\tau\tau} \widetilde{\gamma}_{\tau\tau\tau} J dz' \\
& \leq \varepsilon \|\widehat{\gamma}_{0zzzz}\|_{0, \widehat{\Omega}}^2 + C_1 (\|\widehat{\gamma}\|_{0, \widehat{\Omega}}^2 + \|\widehat{\gamma}_{0z}\|_{2, \widehat{\Omega}}^2 + \|\widetilde{\Gamma}_1\|_{4, \widehat{\Omega}}^2) \\
& \quad + C_2 \left[\|\widehat{\gamma}\|_{4, \widehat{\Omega}}^2 (\|\widehat{\gamma}_0\|_{3, \widehat{\Omega}}^2 + \|\widehat{\gamma}\|_{3, \widehat{\Omega}}^2 + \|\widehat{\gamma}_\sigma\|_{3, \widehat{\Omega}}^2) + \|\widehat{\gamma}\|_{3, \widehat{\Omega}}^2 \left\| \int_0^t \widehat{u} dt' \right\|_{3, \widehat{\Omega}}^2 \right].
\end{aligned}$$

For the quantities

$$\begin{aligned}
& \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \frac{p_{\sigma\widehat{\eta}}}{\widehat{\eta}} \widetilde{\eta}_{\Omega_t n\tau\tau}^2 J dz + c_0 \|\widetilde{\eta}_{\Omega_t n\tau\tau}\|_{0, \widehat{\Omega}}^2, \\
& \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \widehat{\eta} \widetilde{u}_{3n\tau\tau}^2 J dz + c_0 \|\widetilde{u}_{3nn\tau\tau}\|_{0, \widehat{\Omega}}^2, \\
& \|\widetilde{u}'_{z\tau\tau\tau}\|_{0, \widehat{\Omega}}^2 + \|\widetilde{\eta}_{\Omega_t \tau\tau\tau}\|_{0, \widehat{\Omega}}^2, \quad \|\widetilde{u}'_{nn\tau\tau}\|_{0, \widehat{\Omega}}^2, \\
& \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \frac{\widehat{\eta} c_v}{\widehat{\Gamma}} \widetilde{\gamma}_{n\tau\tau}^2 J dz + \frac{\kappa}{\theta^*} \int_{\widehat{\Omega}} \widetilde{\gamma}_{nn\tau\tau}^2 J dz, \\
& \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \frac{p_{\sigma\widehat{\eta}}}{\widehat{\eta}} \widetilde{\eta}_{\Omega_t nn\tau}^2 J dz + \|\widetilde{\eta}_{\Omega_t nn\tau}\|_{0, \widehat{\Omega}}^2, \\
& \|\widetilde{u}_{nnn\tau}\|_{0, \widehat{\Omega}}^2, \quad \|\widetilde{\gamma}_{nnn\tau}\|_{0, \widehat{\Omega}}^2, \quad \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \widehat{\eta} \widetilde{u}_{zzz}^2 J dz, \\
& \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \frac{\widehat{\eta} c_v}{\widehat{\Gamma}} \widetilde{\gamma}_{zzz}^2 J dz
\end{aligned}$$

we obtain the same estimates as in the proof of Lemma 3.9 of [16]. Therefore, from the above considerations we get

$$\begin{aligned}
(3.42) \quad & \frac{1}{2} \frac{d}{dt} \int_{\widehat{\Omega}} \left(\widehat{\eta} \widetilde{u}_{zzz}^2 + \frac{p_{\sigma\widehat{\eta}}}{\widehat{\eta}} \widetilde{\eta}_{\sigma zzz} + \frac{\widehat{\eta} c_v}{\widehat{\Gamma}} \widetilde{\gamma}_{zzz}^2 \right) J dz \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{\widehat{S}} g^{\alpha\beta} \widehat{n} \cdot \int_0^t \widetilde{u}_{\tau\tau\tau s^\alpha} dt' \widehat{n} \cdot \int_0^t \widetilde{u}_{\tau\tau\tau s^\beta} dt' J dz'
\end{aligned}$$

$$\begin{aligned}
& + \frac{\sigma}{2} \frac{d}{dt} \int_{\widehat{S}} (\widehat{H}(\cdot, 0) + 2/R_e)_{,\tau\tau} \widehat{\zeta} \widehat{n} \cdot \int_0^t \widetilde{u}_{\tau\tau\tau_i\tau_i} dt' J dz' \\
& + \frac{\mu}{2} \|\widetilde{u}_{zzzz}\|_{1,\widehat{\Omega}}^2 + \frac{\kappa}{\theta^*} \|\widetilde{\gamma}_{zzzz}\|_{0,\widehat{\Omega}}^2 + c_0 \|\widetilde{\eta}_{\Omega_t zzzz}\|_{0,\widehat{\Omega}}^2 \\
& \leq (\varepsilon + cd)(\|\widetilde{u}_{zzzz}\|_{0,\widehat{\Omega}}^2 + \|\widetilde{\eta}_{\Omega_t zzzz}\|_{0,\widehat{\Omega}}^2 + \|\widetilde{\gamma}_{zzzz}\|_{0,\widehat{\Omega}}^2) \\
& + \varepsilon \left(\|\widetilde{u}_{zzzt}\|_{0,\widehat{\Omega}}^2 + \|\widetilde{\gamma}_{zzzt}\|_{0,\widehat{\Omega}}^2 + \left\| \int_0^t \widetilde{u} dt' \right\|_{4,\widehat{S}}^2 \right. \\
& \quad \left. + \|H(\cdot, 0) + 2/R_e\|_{2,S^1}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{4,S^1}^2 \right) \\
& + C_1 \left(\|\widehat{u}\|_{3,2,\widehat{\Omega}}^2 + \left\| \int_0^t \widehat{u} dt' \right\|_{0,\widehat{\Omega}}^2 + \|\widehat{\gamma}_{0z}\|_{2,\widehat{\Omega}}^2 + \|\widehat{\gamma}\|_{0,\widehat{\Omega}}^2 \right. \\
& \quad \left. + \|\widehat{\eta}_{\sigma z}\|_{1,\widehat{\Omega}}^2 + \|\widehat{\eta}_{\Omega_t}\|_{0,\widehat{\Omega}}^2 + \|\vartheta_{0t}\|_{0,\Omega_t}^2 + \|v\|_{1,\Omega_t}^2 \right. \\
& \quad \left. + \|R(\cdot, t) - R(\cdot, 0)\|_{0,S^1}^2 + \|\widetilde{g}\|_{2,\widehat{\Omega}}^2 + \|\widetilde{k}\|_{2,\widehat{\Omega}}^2 + \|\widetilde{\Gamma}_1\|_{4,\widehat{\Omega}}^2 \right) \\
& + C_2 \left[\left(X_8(\widehat{\Omega}) + \int_0^t \|\widehat{u}\|_{3,\widehat{\Omega}}^2 dt' \right) (1 + X_8^2(\widehat{\Omega})) Y_8(\widehat{\Omega}) \right. \\
& \quad \left. + (\|\widehat{\eta}_\sigma\|_{3,\widehat{\Omega}}^2 + \|\widehat{\gamma}_0\|_{3,\widehat{\Omega}}^2 + \|\widehat{\eta}_\sigma\|_{3,\widehat{\Omega}}^4 + \|\widehat{\eta}_\sigma\|_{3,\widehat{\Omega}}^2 \|\widehat{\gamma}_0\|_{3,\widehat{\Omega}}^2) \right. \\
& \quad \times (\|\vartheta_{0t}\|_{1,\Omega_t}^2 + \|v\|_{2,\Omega_t}^2) \\
& \quad \left. + \|\widehat{u}\|_{3,\widehat{\Omega}}^2 \left\| \int_0^t \widehat{u} dt' \right\|_{4,\widehat{S}}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{3,S^1}^2 \left\| \int_0^t \widehat{u} dt' \right\|_{4,\widehat{S}}^2 \right. \\
& \quad \left. + \|R(\cdot, t) - R(\cdot, 0)\|_{4+1/2,S^1}^2 \left\| \int_0^t \widehat{u} dt' \right\|_{3,\widehat{S}}^2 \right].
\end{aligned}$$

Since the sum of the second and third terms on the left-hand side of (3.42) is equal to

$$\begin{aligned}
& \frac{\sigma}{2} \frac{d}{dt} \int_{\widehat{S}} \frac{1}{2} \widetilde{\delta}^{\alpha\beta} \overline{n} \cdot \int_0^t \widetilde{u}_{p^\gamma p^\delta p^\eta p^\alpha} dt' \overline{n} \cdot \int_0^t \widetilde{u}_{p^\gamma p^\delta p^\eta p^\beta} dt' J dz' \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{\widehat{S}} \left| \widehat{n} \cdot \int_0^t \widetilde{u}_{ppp^1 p^2} dt' \right|^2 J dz' \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{\widehat{S}} \sum_{i=1}^2 \left(\frac{1}{2} \overline{n} \cdot \int_0^t \widetilde{u}_{ppp^i p^i} dt' + 2((\widehat{H}(\cdot, 0) + 2/R_e)\zeta)_{,pp} \right)^2 J dz' \\
& - 4\sigma \frac{d}{dt} \int_{\widehat{S}} (\widehat{H}(\cdot, 0) + 2/R_e)_{,pp}^2 J dz',
\end{aligned}$$

going back to the variables ξ in (3.42), next summing up the resulting inequality and (3.41) over all neighbourhoods of the partition of unity and finally going back to the variables x we get

$$\begin{aligned}
(3.43) \quad & \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_{xxx}^2 + \frac{p_{\sigma\varrho}}{\varrho} \varrho_{xxx}^2 + \frac{\varrho c_v}{\theta} \vartheta_{0xxx}^2 \right) dx \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \frac{1}{2} \tilde{\delta}^{\alpha\beta} \bar{n} \cdot \int_0^t v_{p^\gamma p^\delta p^\eta p^\alpha} dt' \bar{n} \cdot \int_0^t v_{p^\gamma p^\delta p^\eta p^\beta} dt' ds \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \left| \bar{n} \cdot \int_0^t v_{ppp^1 p^2} dt' \right|^2 ds \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} \sum_{i=1}^2 \left(\frac{1}{2} \bar{n} \cdot \int_0^t v_{ppp^i p^i} dt' + 2(H(\cdot, 0) + 2/R_e)_{,pp} \right)^2 ds \\
& \leq \varepsilon \left(\|v_{xxxt}\|_{0,\Omega_t}^2 + \|\vartheta_{0xxxt}\|_{0,\Omega_t}^2 + \left\| \int_0^t v dt' \right\|_{4,S_t}^2 \right. \\
& \quad \left. + \|H(\cdot, 0) + 2/R_e\|_{2,S^1}^2 \|R(\cdot, t) - R(\cdot, 0)\|_{4,S^1}^2 \right) \\
& \quad + C_1 \left(|v|_{3,2,\Omega_t}^2 + \left\| \int_0^t v dt' \right\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{1,\Omega_t}^2 \right. \\
& \quad \left. + \|\vartheta\|_{0,\Omega_t}^2 + \|\vartheta_{0x}\|_{2,\Omega_t}^2 + \|\vartheta_{0t}\|_{2,\Omega_t}^2 \right. \\
& \quad \left. + \|R(\cdot, t) - R(\cdot, 0)\|_{0,S^1}^2 + \|f\|_{2,\Omega_t}^2 + \|r\|_{2,\Omega_t}^2 + \|\theta_1\|_{4,\Omega_t}^2 \right) \\
& \quad + C_2 \left[X_8(1 + X_8^2)Y_8 + \|v\|_{3,\Omega_t}^2 \left\| \int_0^t v dt' \right\|_{4,S_t}^2 \right. \\
& \quad \left. + \|R(\cdot, t) - R(\cdot, 0)\|_{3,S^1}^2 \left\| \int_0^t v dt' \right\|_{4,S_t}^2 \right. \\
& \quad \left. + \|R(\cdot, t) - R(\cdot, 0)\|_{4+1/2,S^1}^2 \left\| \int_0^t v dt' \right\|_{3,S_t}^2 \right. \\
& \quad \left. + \left| \frac{d}{dt} \int_{S_t} (H(\cdot, 0) + 2/R_e)_{,pp}^2 ds \right| \right].
\end{aligned}$$

In view of the interpolation inequality (2.2) we have

$$(3.44) \quad \left| \frac{d}{dt} \int_{S_t} (H(\cdot, 0) + 2/R_e)_{,pp}^2 ds \right| \leq \varepsilon \|v\|_{3,\Omega_t}^2 + C_2 \|H(\cdot, 0) + 2/R_e\|_{2,S^1}^4.$$

Using (3.38) yields

$$\begin{aligned}
(3.45) \quad & \left\| \int_0^t v dt' \right\|_{4,S_t}^2 \leq \varepsilon (\|v_{xxxx}\|_{0,\Omega_t}^2 + \|\varrho_{\sigma xxx}\|_{0,\Omega_t}^2 + \|\vartheta_{0xxx}\|_{0,\Omega_t}^2) \\
& + C_1 \left(\|v\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 + \left\| \int_0^t v dt' \right\|_{0,\Omega_t}^2 \right. \\
& \quad \left. + \|H(\cdot, 0) + 2/R_e\|_{2,S^1}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{4,S^1}^2 \right) \\
& + C_2 (\|v\|_{3,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{2,\Omega_t}^2) \left\| \int_0^t v dt' \right\|_{4,\Omega_t}^2 \\
& \times \left(1 + \left\| \int_0^t v dt' \right\|_{3,\Omega_t}^2 \right).
\end{aligned}$$

From (3.43)–(3.45) we obtain (3.40). ■

The proofs of Lemmas 3.10–3.12 (formulated below) are similar to the proofs of Lemmas 3.10–3.12 of [16]. To estimate the boundary terms associated with the boundary condition (3.1)₄ we use the arguments from Lemmas 4.10–4.12 of [19].

LEMMA 3.10. *Let v, ϱ, ϑ_0 be a sufficiently smooth solution of problem (3.1). Then*

$$\begin{aligned}
& \frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_{xxt}^2 + \frac{p_{\sigma \varrho}}{\varrho} \varrho_{\sigma xxt}^2 + \frac{\varrho c_v}{\theta} \vartheta_{0xxt}^2 \right) dx \\
& + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} g^{\alpha\beta} v_{p^\gamma p^\delta p^\alpha} \cdot \bar{n} v_{p^\gamma p^\delta p^\beta} \cdot \bar{n} ds \\
& + c_0 (\|v_{xxt}\|_{1,\Omega_t}^2 + \|\varrho_{\sigma xxt}\|_{0,\Omega_t}^2 + \|\vartheta_{0xxt}\|_{0,\Omega_t}^2) \\
& \leq \varepsilon (\|v_{xxtt}\|_{0,\Omega_t}^2 + \|v_{xxxx}\|_{0,\Omega_t}^2 + \|\vartheta_{0xxxx}\|_{0,\Omega_t}^2 \\
& \quad + \|\vartheta_{0xxtt}\|_{0,\Omega_t}^2 + \|H(\cdot, 0) + 2/R_e\|_{0,S^1}^2) \\
& \quad + C_1 (\|v\|_{3,1,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{1,\Omega_t}^2 + \|\varrho_{\sigma t}\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 \\
& \quad + \|\vartheta_{0x}\|_{2,\Omega_t}^2 + \|\vartheta_{0t}\|_{2,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 + |f|_{2,1,\Omega_t}^2 \\
& \quad + |r|_{2,1,\Omega_t}^2 + \|\theta_{1t}\|_{3,\Omega_t}^2 + \|\theta_1\|_{3,\Omega_t}^2) + C_2 X_9 (1 + X_9^2) Y_9,
\end{aligned}$$

where the summation over repeated indices ($\alpha, \beta, \gamma, \delta = 1, 2$) and coordinates x is assumed and

$$\begin{aligned} X_9 &= |v|_{3,2,\Omega_t}^2 + |\varrho_\sigma|_{3,1,\Omega_t}^2 + |\vartheta_0|_{3,1,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{3,\Omega_{t'}}^2 dt', \\ Y_9 &= |v|_{4,3,\Omega_t}^2 + |\varrho_\sigma|_{3,1,\Omega_t}^2 + \|\vartheta_{0x}\|_{3,\Omega_t}^2 + \|\vartheta_{0t}\|_{3,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 \\ &\quad + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{4,\Omega_{t'}}^2 dt'. \end{aligned}$$

LEMMA 3.11. Let v, ϱ, ϑ_0 be a sufficiently smooth solution of problem (3.1). Then

$$\begin{aligned} &\frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_{xtt}^2 + \frac{p_{\sigma\varrho}}{\varrho} \varrho_{\sigma xtt}^2 + \frac{\varrho c_v}{\theta} \vartheta_{0xtt}^2 \right) dx + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} g^{\alpha\beta} \bar{n} \cdot v_{tpp^\alpha} \bar{n} \cdot v_{tpp^\beta} ds \\ &\quad + c_0 (\|v_{txx}\|_{1,\Omega_t}^2 + \|\varrho_{\sigma tt}\|_{1,\Omega_t}^2 + \|\vartheta_{0txx}\|_{0,\Omega_t}^2) \\ &\leq \varepsilon (\|v_{xxxt}\|_{0,\Omega_t}^2 + \|v_{xttt}\|_{0,\Omega_t}^2 + \|\vartheta_{0xxx}\|_{0,\Omega_t}^2 \\ &\quad + \|\vartheta_{0xttt}\|_{0,\Omega_t}^2 + \|H(\cdot, 0) + 2/R_e\|_{2,S^1}^2) \\ &\quad + C_1 (|v|_{3,0,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{1,\Omega_t}^2 + \|\varrho_{\sigma t}\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 \\ &\quad + \|\vartheta_{0x}\|_{2,\Omega_t}^2 + \|\vartheta_{0t}\|_{2,0,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 + |f|_{2,0,\Omega_t}^2 \\ &\quad + |r|_{2,0,\Omega_t}^2 + \|\theta_{1tt}\|_{2,\Omega_t}^2 + \|\theta_{1t}\|_{2,\Omega_t}^2 + \|\theta_1\|_{2,\Omega_t}^2) \\ &\quad + C_2 X_{10} (1 + X_{10}^2) Y_{10}, \end{aligned}$$

where

$$\begin{aligned} X_{10} &= |v|_{3,0,\Omega_t}^2 + |\varrho_\sigma|_{3,0,\Omega_t}^2 + |\vartheta_0|_{3,0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{3,\Omega_{t'}}^2 dt', \\ Y_{10} &= |v|_{4,1,\Omega_t}^2 + |\varrho_\sigma|_{3,1,\Omega_t}^2 + \|\vartheta_{0x}\|_{3,\Omega_t}^2 + \|\vartheta_{0t}\|_{3,1,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2 \\ &\quad + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{4,\Omega_{t'}}^2 dt'. \blacksquare \end{aligned}$$

LEMMA 3.12. Let v, ϱ, ϑ_0 be a sufficiently smooth solution of problem (3.1). Then

$$\begin{aligned} &\frac{1}{2} \frac{d}{dt} \int_{\Omega_t} \left(\varrho v_{ttt}^2 + \frac{p_{\sigma\varrho}}{\varrho} \varrho_{\sigma ttt}^2 + \frac{\varrho c_v}{\theta} \vartheta_{0ttt}^2 \right) dx + \frac{\sigma}{2} \frac{d}{dt} \int_{S_t} g^{\alpha\beta} \bar{n} \cdot v_{tts^\alpha} \bar{n} \cdot v_{tts^\beta} ds \\ &\quad + c_0 (\|v_{ttt}\|_{1,\Omega_t}^2 + \|\varrho_{\sigma ttt}\|_{0,\Omega_t}^2 + \|\vartheta_{0ttt}\|_{1,\Omega_t}^2) \\ &\leq C_1 (\|v_{tt}\|_{1,\Omega_t}^2 + \|\vartheta_{0tt}\|_{1,\Omega_t}^2 + \|f_{ttt}\|_{0,\Omega_t}^2 + |f|_{2,0,\Omega_t}^2 \\ &\quad + \|r_{ttt}\|_{0,\Omega_t}^2 + |r|_{2,0,\Omega_t}^2 + |\theta_1|_{4,1,\Omega_t}^2) + C_2 X_{11} (1 + X_{11}^3) Y_{11}, \end{aligned}$$

where

$$\begin{aligned} X_{11} &= |v|_{3,0,\Omega_t}^2 + |\varrho_\sigma|_{3,0,\Omega_t}^2 + |\vartheta_0|_{3,0,\Omega_t}^2, \\ Y_{11} &= |v|_{4,1,\Omega_t}^2 + |\varrho_\sigma|_{3,0,\Omega_t}^2 + |\vartheta_0|_{4,1,\Omega_t}^2. \blacksquare \end{aligned}$$

Estimating $\|\vartheta\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2$ by $\|\vartheta_{0x}\|_{0,\Omega_t}^2 + \|p_\sigma\|_{0,\Omega_t}^2$ (by using (3.20) and (3.21)) from the above lemmas for sufficiently small ε we obtain

THEOREM 3.13. *Let*

$$\begin{aligned} \phi(t) = & \int_{\Omega_t} \varrho \sum_{0 \leq |\alpha|+i \leq 3} |D_x^\alpha \partial_t^i v|^2 dx + \int_{\Omega_t} \left(\frac{p_1}{\varrho} \varrho_\sigma^2 + \bar{\varrho}_{\Omega_t}^2 + \frac{p_2 \varrho c_v}{p_\theta \theta} \vartheta_0^2 \right) dx \\ & + \int_{\Omega_t} \frac{p_\sigma \varrho}{\varrho} \sum_{1 \leq |\alpha|+i \leq 3} |D_x^\alpha \partial_t^i \varrho_\sigma|^2 dx + \int_{\Omega_t} \frac{\varrho c_v}{\theta} \sum_{1 \leq |\alpha|+i \leq 3} |D_x^\alpha \partial_t^i \vartheta_0|^2 dx \\ & + \frac{\sigma}{2} \int_{S_t} \tilde{\delta}^{\alpha\beta} \sum_{|k| \leq 2} \bar{n} \cdot \int_0^t \partial_p^k v_{p^\gamma p^\alpha} dt' \bar{n} \cdot \int_0^t \partial_p^k v_{p^\gamma p^\beta} dt' ds \\ & + \frac{\sigma}{2} \int_{S_t} \sum_{|k| \leq 2} \left| \bar{n} \cdot \int_0^t \partial_p^k v_{p^1 p^2} dt' \right|^2 ds \\ & + \frac{\sigma}{2} \int_{S_t} \sum_{|k| \leq 2} \sum_{i=1}^2 \left(\frac{1}{2} \bar{n} \cdot \int_0^t \partial_p^k v_{p^i p^i} dt' + 2 \partial_p^k (H(\cdot, 0) + 2/R_e) \right)^2 ds \\ & + \frac{\sigma}{2} \int_{S_t} g^{\alpha\beta} \left(\bar{n} \cdot \int_0^t v_{s^\alpha} dt' \bar{n} \cdot \int_0^t v_{s^\beta} dt' + \sum_{|k| \leq 2} D_t^k v_{s^\alpha} \cdot \bar{n} D_t^k v_{s^\beta} \cdot \bar{n} \right. \\ & \quad \left. + \sum_{|k| \leq 2} D_p^k v_{p^\alpha} \cdot \bar{n} D_p^k v_{p^\beta} \cdot \bar{n} \right) ds \end{aligned}$$

($\tilde{\delta}^{\alpha\beta}$ is defined in Lemma 3.4),

$$\begin{aligned} \Phi(t) = & |v|_{4,1,\Omega_t}^2 + |\varrho_\sigma|_{3,0,\Omega_t}^2 - \|\varrho_\sigma\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 \\ & + |\vartheta_0|_{4,1,\Omega_t}^2 - \|\vartheta_0\|_{0,\Omega_t}^2 + \|\vartheta\|_{0,\Omega_t}^2, \\ \psi(t) = & \|v\|_{0,\Omega_t}^2 + \|p_\sigma\|_{0,\Omega_t}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{0,S^1}^2, \\ F(t) = & \|f_{ttt}\|_{0,\Omega_t}^2 + |f|_{2,0,\Omega_t}^2 + \|r_{ttt}\|_{0,\Omega_t}^2 + |r|_{2,0,\Omega_t}^2 \\ & + \|r\|_{0,\Omega_t} + |\theta_1|_{4,1,\Omega_t}^2 + \|\theta_1\|_{1,\Omega_t}. \end{aligned}$$

Assume that $\nu > \frac{1}{3}\mu$. Then for sufficiently smooth solutions of problem (3.1) the following estimate holds:

$$\begin{aligned} (3.46) \quad & \frac{d\phi}{dt} + c_0 \Phi \leq c_1 P(X) X (1 + X^3) (X + Y) + c_2 F \\ & + c_3 \psi + c_4 \|H(\cdot, 0) + 2/R_e\|_{2,S^1}^4 \\ & + \varepsilon c_5 (\|H(\cdot, 0) + 2/R_e\|_{2,S^1}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{4,S^1}^2) \\ & + c_6 \left(\|R(\cdot, t) - R(\cdot, 0)\|_{4+1/2,S^1}^2 \left\| \int_0^t v dt' \right\|_{3,S_t}^2 \right. \\ & \quad \left. + \|r\|_{0,\Omega_t}^2 + |\theta_1|_{4,1,\Omega_t}^2 + \|\theta_1\|_{1,\Omega_t} \right) \end{aligned}$$

$$+ \|R(\cdot, t) - R(\cdot, 0)\|_{3,S^1}^2 \left\| \int_0^t v dt' \right\|_{4,S_t}^2,$$

where $0 < c_0 < 1$ is a constant depending on ϱ_* , ϱ^* , θ_* , θ^* , μ , ν and κ , c_i ($i=1, \dots, 6$) depend on ϱ_* , ϱ^* , θ_* , θ^* , T , $\int_0^T \|v\|_{3,\Omega_{t'}}^2 dt'$, $\|S\|_{4+1/2}$, constants from the imbedding lemma and the Korn inequalities (see Section 5 of [18]), ε is a small parameter and

$$\begin{aligned} X &= |v|_{3,0,\Omega_t}^2 + |\varrho_\sigma|_{3,0,\Omega_t}^2 + |\vartheta_0|_{3,0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{3,\Omega_{t'}}^2 dt', \\ Y &= |v|_{4,1,\Omega_t}^2 + |\varrho_{\sigma t}|_{2,0,\Omega_t}^2 + \|\varrho_{\sigma x}\|_{2,\Omega_t}^2 + |\vartheta_{0t}|_{3,1,\Omega_t}^2 + \|\vartheta_{0x}\|_{3,\Omega_t}^2 \\ &\quad + \|\vartheta\|_{0,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2 + \int_0^t \|v\|_{4,\Omega_{t'}}^2 dt'. \blacksquare \end{aligned}$$

A slight modification of the proof of Theorem 3.13 yields

THEOREM 3.14. *Assume that $\nu > \frac{1}{3}\mu$. For sufficiently smooth solutions of problem (3.1) we have*

$$\begin{aligned} (3.47) \quad \frac{d\tilde{\phi}}{dt} + c_0\Phi &\leq c_7 P \left(\tilde{\phi} + \int_0^t \|v\|_{4,\Omega_{t'}}^2 dt' \right) \left(\tilde{\phi} + \int_0^t \|v\|_{4,\Omega_{t'}}^2 dt' \right) \\ &\quad \times \left[1 + \left(\tilde{\phi} + \int_0^t \|v\|_{4,\Omega_{t'}}^2 dt' \right)^3 \right] \left(\Phi + \tilde{\phi} + \int_0^t \|v\|_{4,\Omega_{t'}}^2 dt' \right) \\ &\quad + c_8 F + c_9 \psi + c_{10} \|H(\cdot, 0) + 2/R_e\|_{2,S^1}^4 \\ &\quad + \varepsilon c_{11} (\|H(\cdot, 0) + 2/R_e\|_{2,S^1}^2 + \|R(\cdot, t) - R(\cdot, 0)\|_{4,S^1}^2) \\ &\quad + c_{12} \left(\|R(\cdot, t) - R(\cdot, 0)\|_{4-1/2,S^1}^2 \left\| \int_0^t v dt' \right\|_{3,S_t}^2 \right. \\ &\quad \left. + \|R(\cdot, t) - R(\cdot, 0)\|_{3,S^1}^2 \left\| \int_0^t v dt' \right\|_{4,S_t}^2 \right), \end{aligned}$$

where ε is a small parameter, c_0 and c_i ($i=7, \dots, 12$) have the same properties as in Theorem 3.13 and $\tilde{\phi}(t) = |v|_{3,2,\Omega_t}^2 + |\varrho_\sigma|_{3,2,\Omega_t}^2 + |\vartheta_0|_{3,2,\Omega_t}^2 + \|\bar{\varrho}_{\Omega_t}\|_{0,\Omega_t}^2$.

References

- [1] O. V. Besov, V. P. Il'in and S. M. Nikol'skii, *Integral Representations of Functions and Imbedding Theorems*, Nauka, Moscow, 1975 (in Russian).
- [2] L. Landau and E. Lifschitz, *Mechanics of Continuum Media*, Nauka, Moscow, 1984; new edition: *Hydrodynamics*, Nauka, Moscow, 1986 (in Russian).

- [3] A. Matsumura and T. Nishida, *The initial value problem for the equations of motion of viscous and heat-conductive gases*, J. Math. Kyoto Univ. 20 (1980), 67–104.
- [4] —, —, *The initial value problem for the equations of motion of compressible viscous and heat-conductive fluids*, Proc. Japan Acad. Ser. A 55 (1979), 337–342.
- [5] —, —, *The initial boundary value problem for the equations of motion of compressible viscous and heat-conductive fluid*, preprint Univ. of Wisconsin, MRC Technical Summary Report 2237 (1981).
- [6] —, —, *Initial boundary value problems for the equations of motion of general fluids*, in: Computing Methods in Applied Sciences and Engineering, R. Glowinski and J. L. Lions (eds.), North-Holland, Amsterdam, 1982, 389–406.
- [7] —, —, *Initial boundary value problems for the equations of motion of compressible viscous and heat-conductive fluids*, Comm. Math. Phys. 89 (1983), 445–464.
- [8] K. Pileckas and W. M. Zajączkowski, *On free boundary problem for stationary compressible Navier–Stokes equations*, ibid. 128 (1990), 1–36.
- [9] V. A. Solonnikov and A. Tani, *Evolution free boundary problem for equations of motion of viscous compressible barotropic liquid*, preprint of Paderborn University.
- [10] A. Valli, *Periodic and stationary solutions for compressible Navier–Stokes equations via a stability method*, Ann. Scuola Norm. Sup. Pisa (4) 10 (1983), 607–647.
- [11] A. Valli and W. M. Zajączkowski, *Navier–Stokes equations for compressible fluids: global existence and qualitative properties of the solutions in the general case*, Comm. Math. Phys. 103 (1986), 259–296.
- [12] E. Zadrzyńska and W. M. Zajączkowski, *On local motion of a general compressible viscous heat conducting fluid bounded by a free surface*, Ann. Polon. Math. 59 (1994), 133–170.
- [13] —, —, *On global motion of a compressible heat conducting fluid bounded by a free surface*, Acta Appl. Math. 37 (1994), 221–231.
- [14] —, —, *Conservation laws in free boundary problems for viscous compressible heat conducting fluids*, Bull. Polish Acad. Sci. Tech. Sci. 42 (1994), 195–205.
- [15] —, —, *Conservation laws in free boundary problems for viscous compressible heat conducting capillary fluids*, ibid. 43 (1995), 423–444.
- [16] —, —, *On a differential inequality for a viscous compressible heat conducting fluid bounded by a free surface*, Ann. Polon. Math. 61 (1995), 141–188.
- [17] —, —, *On the global existence theorem for a free boundary problem for equations of a viscous compressible heat conducting fluid*, ibid. 63 (1996), 199–221.
- [18] W. M. Zajączkowski, *On nonstationary motion of a compressible barotropic viscous fluid bounded by a free surface*, Dissertationes Math. 324 (1993).
- [19] —, *On nonstationary motion of a compressible barotropic viscous capillary fluid bounded by a free surface*, SIAM J. Math. Anal. 25 (1994), 1–84.

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