

Application of the extremal points method to some variational problems in the theory of schlicht functions

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The aim of this paper is to show the application of Leja's method to some variational problems in the theory of schlicht functions. Let $z = \varphi(w)$, $\varphi(\infty) = \infty$ be the conformal mapping of the exterior of the unit circle $|w| > 1$ onto the unbounded domain D which contains the point $z = \infty$. The complementary set E is a bounded continuum. Suppose that the capacity $d(E)$ of E is 1. Denote by $w = \Phi(z)$ the inverse function to $z = \varphi(w)$. We have

$$(1) \quad \log \Phi(z) = G(z, \infty) + iH(z, \infty),$$

where $G(z, \infty)$ is the Green function of D with the pole at $z = \infty$, $H(z, \infty)$ is its conjugate.

If E^* denotes a new continuum obtained from E by a variation, we shall later specify the kind of variation, and Φ^* , G^* , H^* the corresponding functions to Φ , G , H defined above, then

$$(2) \quad \log \Phi^*(z) = G^*(z, \infty) + iH^*(z, \infty).$$

Let J be a Jordan curve which contains E strictly inside, i.e. the distance from J to E is positive. For small enough variation of E the continuum E^* lies inside J . If we put

$$(3) \quad h(z) = G^*(z, \infty) - G(z, \infty) + i(H^*(z, \infty) - H(z, \infty))$$

then for $|h(z)| < \varepsilon$, $z \in J$ is $|h(z)| < \varepsilon$ for z outside of J .

From (1), (2) and (3) follows $\Phi^*(z) = \Phi(z)e^{h(z)}$. Therefore

$$(4) \quad z = \Phi^{*-1}(\Phi(z)e^{h(z)}) = \Phi^{*-1}(w + wh + O(h^2))$$

where $O(h^2)$ denotes the terms which contain h in higher powers.

If one extends the right side of (4) with respect to h one obtains ([1], p. 302)

$$(5) \quad z = z^* + whz' + O(h^2).$$

Under the transformation $w = 1/x$, $z = 1/y$ the formula (5) takes the form

$$(6) \quad y^* = y + xhy' + O(h^2).$$

Here $y = y(x)$ denotes the conformal mapping of the unit circle $|x| < 1$ onto Δ , where Δ is the image of D under $y = 1/z$. Now we consider the variation of D defined by

$$(7) \quad z^* = z/(1 + z\rho e^{ia}),$$

where $\rho \geq 0$ and a is an arbitrary real number. Let D^* be the image of D under (7). For $\rho = 0$ is $D^* = D$; for $\rho > 0$ and $|a|$ small enough D^* represents small variation of D .

According to (7) is $(z = \infty) \Leftrightarrow (z^* = 1/\rho e^{ia})$ and $(z^* = \infty) \Leftrightarrow (z = -1/\rho e^{ia})$. We want to compute (3) in the case of variation (7). Observe that

$$\zeta = \exp[G(z; -1/\rho e^{ia}) + iH(z; -1/\rho e^{ia})]$$

and

$$\zeta^* = \exp[G(z; \infty) + iH(z; \infty)]$$

are related by the homography

$$\begin{aligned} \zeta &= \frac{1 - \bar{\zeta}_0^* \zeta^*}{\zeta^* - \bar{\zeta}_0^*} e^{i\theta}, \quad \zeta_0^* = \exp\left[G\left(-\frac{1}{\rho e^{ia}}; \infty\right) + iH\left(-\frac{1}{\rho e^{ia}}; \infty\right)\right], \\ &\exp\left\{i\theta + G\left(-\frac{1}{\rho e^{ia}}; \infty\right) + iH\left(-\frac{1}{\rho e^{ia}}; \infty\right)\right\} \\ &= -\exp\left\{G\left(\infty; -\frac{1}{\rho e^{ia}}\right) + iH\left(\infty; -\frac{1}{\rho e^{ia}}\right)\right\}, \\ \theta &= 2H\left(-\frac{1}{\rho e^{ia}}; \infty\right). \end{aligned}$$

Therefore

$$(8) \quad \exp\left\{G\left(z; -\frac{1}{\rho e^{ia}}\right) + iH\left(z; -\frac{1}{\rho e^{ia}}\right)\right\} \\ = e^{i\theta} \frac{1 - \exp\left[G\left(-\frac{1}{\rho e^{ia}}; \infty\right) - iH\left(-\frac{1}{\rho e^{ia}}; \infty\right) + G(z; \infty) + iH(z, \infty)\right]}{\exp[G(z, \infty) + iH(z, \infty)] - \exp\left[G\left(-\frac{1}{\rho e^{ia}}; \infty\right) + iH\left(-\frac{1}{\rho e^{ia}}; \infty\right)\right]}$$

and

$$(9) \quad h(z) = G\left(\eta; -\frac{1}{\rho e^{ia}}\right) + iH\left(\eta; -\frac{1}{\rho e^{ia}}\right) - G(z, \infty) - iH(z, \infty),$$

where

$$(10) \quad \eta = z/(1 - z\rho e^{ia}).$$

From (8) and (9) follows for small $|h|$

$$1 + h \approx \exp h$$

$$= e^{i\theta} \frac{f(z)}{\exp[G(\eta; \infty) + iH(\eta; \infty)] - \exp\left[G\left(-\frac{1}{\rho e^{ia}}; \infty\right) + iH\left(-\frac{1}{\rho e^{ia}}; \infty\right)\right]},$$

where

$$f(z) = \exp[-G(z, \infty) - iH(z, \infty)] - \exp\left[G\left(-\frac{1}{\rho e^{ia}}; \infty\right) - iH\left(-\frac{1}{\rho e^{ia}}; \infty\right) + G(\eta; \infty) + iH(\eta; \infty) - G(z, \infty) - iH(z, \infty)\right].$$

For small $\rho > 0$, we neglect the terms with ρ in higher power than 1, is

$$\exp\left[-G\left(-\frac{1}{\rho e^{ia}}; \infty\right) - iH\left(-\frac{1}{\rho e^{ia}}; \infty\right)\right] \approx -\rho e^{ia}.$$

Therefore using the fact that (see [2])

$$G(z; \infty) + iH(z; \infty) = \int \log(z - \xi) d\mu(\xi),$$

where μ is the natural mass-distribution on E , we have

$$1 + h(z) \approx \frac{1 + \int \log(\eta - \xi) d\mu(\xi) - \int \log(z - \xi) d\mu(\xi) + \rho e^{ia} \exp[-\int \log(z - \xi) d\mu]}{1 + \rho e^{ia} \exp[\int \log(\eta - \xi) d\mu(\xi)]}$$

and

$$(11) \quad h(z) \approx \int \log(\eta - \xi) d\mu(\xi) - \int \log(z - \xi) d\mu(\xi) + \rho e^{ia} \exp[-\int \log(z - \xi) d\mu] - \rho e^{ia} \exp[\int \log(\eta - \xi) d\mu].$$

On the other hand, for $\xi \in E$ and $z \in D$ such that $|\xi/z| < 1$ is

$$\int \log(z - \xi) d\mu = \log z + \int \log\left(1 - \frac{\xi}{z}\right) d\mu = \log z - \frac{s_1}{z} - \frac{s_2}{2z^2} - \frac{s_3}{3z^3} - \dots,$$

where $s_k = \int \xi^k d\mu$, see [3], $k = 1, 2, \dots$

From (10) and (11) follows that

$$(12) \quad h(z) = \rho \left\{ e^{ia} \left[\left(z + s_1 + \frac{s_2}{z} + \frac{s_3}{z^2} + \dots \right) - z \left(1 - \frac{s_1}{z} - \frac{s_2}{2z^2} - \frac{s_3}{3z^3} - \dots \right) \right] + e^{-ia} \frac{1}{z} \left(1 + \frac{s_1}{z} + \frac{s_2}{2z^2} + \dots \right) \right\} + O(\rho^2).$$

As the function $y(x)$ belongs to the class S of all schlicht functions in the unit circle with the extension

$$(13) \quad y = x + a_2 x^2 + a_3 x^3 + \dots$$

we can use the variational formula (6) to compare the coefficients of (13) with those of

$$(14) \quad y^* = a_1^* x + a_2^* x^2 + a_3^* x^3 + \dots$$

If we put (13) and (14) in (6), using (12) and compare the coefficients on both sides, we obtain

$$(15) \quad a_1^* = 1 + 2\rho e^{i\alpha} s_1,$$

$$(16) \quad a_2^* = a_2 + \rho \{e^{i\alpha}[-4a_2^2 + \frac{3}{2}s_2 - \frac{1}{2}s_1^2] + e^{-i\alpha}\},$$

$$(17) \quad a_3^* = a_3 + \rho \{e^{i\alpha}[-4a_3] + 2a_2 e^{-i\alpha}\}, \dots$$

The function $y^*(x)/a_1^*$ belongs to S . In [3] are given formulas which permit to express the coefficients a_k by s_k . For example

$$(18) \quad a_2 = -s_1; \quad a_3 = \frac{3s_1^2 - s_2}{2}; \quad a_4 = \frac{-8s_1^3 + 6s_1s_2 - s_3}{3}, \quad \dots$$

Suppose that $y(x)$ is this function of S with $a_2 > 0$ which realise the maximum of $|a_2|$ in S . For this function is

$$(19) \quad |a_2^*/a_1^*| \leq a_2.$$

Using (15) and (16) we obtain from (19)

$$(20) \quad \frac{a_2^*}{a_1^*} = a_2 + \rho \{e^{i\alpha}[-3a_3 + 2a_2^2] + e^{-i\alpha}\},$$

$$|a_2 + \rho \{e^{i\alpha}[-3a_3 + 2a_2^2] + e^{-i\alpha}\}| \leq a_2$$

for every $\rho > 0$ and a . (20) can be rewrite in the form

$$\operatorname{Re} \{e^{i\alpha}(-3a_3 + 2a_2^2 + 1)\} \leq 0.$$

Therefore

$$(21) \quad 3a_3 = 2a_2^2 + 1.$$

The condition (21) is the special case of Marty's formula for $n = 2$ (see [4]). In the similar way we obtain from (15) and (17) for the function of S with $a_3 > 0$ which realises the maximum of $|a_3|$

$$(22) \quad \frac{a_3^*}{a_1^*} = a_3 + \rho \{e^{i\alpha}[-4a_4 + 2a_2a_3] + 2a_2 e^{-i\alpha}\}$$

and in consequence ($a_3 > 0$)

$$(23) \quad 4a_4 = 2a_2a_3 + 2\bar{a}_2.$$

This is Marty's formula for $n = 3$.

If we suppose that the second coefficient of the extremal function has the argument φ_2 , i.e. if $a_2 = |a_2|e^{i\varphi_2}$, then (21) takes the form

$$(21^*) \quad 3a_3 = 2a_2^2 + e^{2i\varphi_2}.$$

In the similar way if the third coefficient a_3 of the extremal function is equal $a_3 = |a_3|e^{i\varphi_3}$, then

$$(22^*) \quad 4a_4 = 2a_2a_3 + 2\bar{a}_2e^{2i\varphi_3}.$$

Denote by b_k the coefficients of the inverse function $x = x(y)$

$$x = y + b_2y^2 + b_3y^3 + \dots$$

Using (20) and (22) we obtain easily the following conditions for the coefficients b_2, b_3, b_4 of the function $x(y)$ which realises the maximum of $|b_2|$ or $|b_3|$ respectively

$$\begin{aligned} 3b_3 &= 4b_2^2 - e^{2i\varphi_2}, & b_2 &= |b_2|e^{i\varphi_2}, \\ 4b_4 &= 6b_2b_3 + 2\bar{b}_2e^{2i\varphi_3}, & b_3 &= |b_3|e^{i\varphi_3}. \end{aligned}$$

References

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