

CONIC DEGENERATION OF THE DIRAC OPERATOR

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0. Introduction. Consider a manifold M with an embedded hypersurface N . We obtain a family of Riemannian manifolds by shrinking N to a point, in such a way that the limiting manifold has a conic singularity with two copies of N as cross section. In this situation one can ask: What happens to the basic geometric operators in this limit? The case of $\bar{\partial}$ on a Riemann surface was treated in [SS]. Here, we consider the Dirac operator; a closely related paper [S] treats the Gauss–Bonnet operator.

The “shrinking” process is the following. Take a fixed metric \tilde{g} on N , and let x , $|x| < 1$, be a coordinate transversal to N . We define a family of metrics g_t on M , depending on the parameter t , by

$$(0.1) \quad g_t = dx^2 + (x^2 + t^2)\tilde{g}, \quad |x| < 1.$$

For $|x| \geq 1$, the metric is extended in some convenient way. In the limit as $t \rightarrow 0$, this is a conic metric; the cross section of the cone is two copies of N . Denote by D_t the full, selfadjoint Dirac operator on M , with the metric g_t ; by D_{\max} the maximal realization of the Dirac operator on the conic limiting manifold; and by \tilde{D} the Dirac operator on N . As shown in [C], and at the end of Sec. 1 below, each eigenvalue μ of \tilde{D} with $|\mu| < 1/2$ contributes sections in the domain of D_{\max} with singularities like $|x|^{-(n/2) \pm \mu}$, which are therefore in L^2 . When such sections are present, D is not essentially selfadjoint, and there are various possible closed realizations (see [C] and [BS]). Our limit picks out a particular one.

THEOREM 1. *The graph of D_t has a limit as $t \rightarrow 0$, which is the graph of a selfadjoint Dirac operator D_0 on the conic limiting manifold. The domain of D_0 consists of the sections in the domain of D_{\max} which satisfy the condition*

$$(0.2) \quad \lim_{x \rightarrow 0^+} |x|^{n/2} s(x) = \lim_{x \rightarrow 0^-} |x|^{n/2} s(x).$$

The limits in (0.2) are taken in $L^2(N)$. The existence of these limits means that the singularities $|x|^{-(n/2) - |\mu|}$ are required to vanish for $\mu \neq 0$; sections where the limits in (0.2) are finite and nonzero arise only from

the harmonic spinors on N , with eigenvalue $\mu = 0$. Thus, when no such harmonic spinors exist, the link between the two sides of N dissolves completely; in the limit, there is no condition relating the sections on one side of N to sections on the other side.

When N is the standard unit sphere, then the limiting metric g_0 is the standard metric on the unit ball; so functions in the maximal domain are in the Sobolev space H^1 , ruling out any singularities $|x|^{-(n/2)\pm|\mu|}$ for $|\mu| < 1/2$. This gives a peculiar proof that the Dirac operator on the standard unit sphere has no eigenvalue with $|\mu| < 1/2$. (In the case of S^1 , we have to specify the spin structure; it is the nontrivial one arising by considering S^1 as the boundary of the unit 2-disk.)

Now suppose that $\dim(M)$ is even, and let D_t^+ denote the Dirac operator D_t restricted to the “+ spinors”. The index of D_t^+ is the \hat{A} genus of M . Now each eigenvalue of \tilde{D} with $|\mu| < 1/2$ contributes on each side of N just one section in the domain of D_{\max} , with singularity $x^{-(n/2)+\mu}$ as $x \rightarrow 0^+$ and $|x|^{-(n/2)-\mu}$ as $x \rightarrow 0^-$ (see end of Sec. 1).

THEOREM 2. *The graph of D_t^+ has a limit as $t \rightarrow 0$, which is the graph of a Dirac operator D_0^+ on the conic limiting manifold. The domain of D_0^+ is again determined by the condition (0.2).*

Since graph continuity preserves the index (see [SS]), we obtain:

COROLLARY 1. *The limiting operator D_0^+ has index $\text{ind}(D_0^+) = \hat{A}(M)$.*

Suppose now that N separates M into two disjoint parts M_+ and M_- . Then we obtain two disjoint Riemannian manifolds \overline{M}_+ and \overline{M}_- , each with a conic singularity with cross section N . On each of these, we can consider Dirac operators with appropriate domains. Denote by $D_{\text{Dir}}^+(\overline{M}_{\pm})$ the restriction of $D_{\max}^+(\overline{M}_{\pm})$ to sections with

$$|x|^{n/2}s(x) = o(1),$$

and by $D_{\text{Neu}}^+(\overline{M}_{\pm})$ the restriction to sections with $o(1)$ replaced by $O(1)$. The distinction between these two realizations lies in the treatment of the harmonic spinors, which give rise to sections where the limits in (0.2) are finite and nonzero; the Dirichlet domain requires these limits to vanish, while the Neumann domain leaves them free. Denote the indices of these realizations by $\hat{A}_{\text{Dir}}(\overline{M}_{\pm})$ and $\hat{A}_{\text{Neu}}(\overline{M}_{\pm})$. Then Theorem 2 yields:

COROLLARY 2. *In the above situation*

$$\hat{A}(M) = \hat{A}_{\text{Dir}}(\overline{M}_+) + \hat{A}_{\text{Neu}}(\overline{M}_-) = \hat{A}_{\text{Dir}}(\overline{M}_+) + \hat{A}_{\text{Dir}}(\overline{M}_-) + \nu(\tilde{D})$$

where $\nu(\tilde{D})$ is the nullity of the Dirac operator on N .

The outline of the paper is as follows. Section 1 reviews the separation of variables formula for the Dirac operator given by Chou [C]. Section 2

constructs a family of parametrices for D_t near the hypersurface N , and studies their limit as $t \rightarrow 0$. Section 3 reviews the argument for graph continuity from [SS], thus proving Theorems 1 and 2. In the process, we include a part of the argument that was regrettably omitted from [SS].

The author thanks Haynes Miller and I. M. Singer for helpful comments about $\widehat{A}(M)$ and the Dirac operator, and J. Bruening for a careful reading of an earlier draft.

1. The representations of D and D^+ using eigenvectors of \widetilde{D} . This section is essentially a review of part of [C].

A metric g on M gives an inner product on $T(M)$, and thus determines a Clifford bundle on M , a vector bundle whose fibre at any point p is the tensor algebra generated by the tangent vectors, reduced modulo the relations

$$(1.1) \quad \partial/\partial x_1 \otimes \partial/\partial x_2 + \partial/\partial x_2 \otimes \partial/\partial x_1 = -2\langle \partial/\partial x_1, \partial/\partial x_2 \rangle_g.$$

We denote by \cdot the multiplication in this Clifford bundle. If M is a spin manifold, it carries a spin bundle S on which the Clifford bundle acts. The metric g defines a connection on S . The Dirac operator is then expressed, using an orthonormal frame $\{e_j\}$ for $T(M)$, as

$$(1.2) \quad D = \sum e_j \cdot \nabla_{e_j}.$$

With our local representation near N , we can choose an orthonormal frame $\{e_1, \dots, e_n\}$ for N , and supplement it with $e_0 = \partial/\partial x$.

When $\dim(M) = n + 1$ is even, then S is a direct sum of subbundles S^\pm determined by the action of the volume element

$$\omega_{n+1} = i^{(n+1)/2} \frac{\partial}{\partial x} \cdot \omega_n, \quad \omega_n = e_1 \cdot \dots \cdot e_n.$$

Since $\omega_{n+1} \cdot \omega_{n+1} = 1$, subbundles S^\pm of S can be (and are) defined by

$$\omega_{n+1} \cdot s^\pm = \pm s^\pm \quad \text{for any section } s^\pm \text{ in } S^\pm.$$

This gives a restriction of D ,

$$D^+ : C^\infty(S^+) \rightarrow C^\infty(S^-).$$

The restriction $S^+|_N$ determines a spin bundle $S(N)$ on N , and as in [C], we denote by \widetilde{D} the Dirac operator on this $S(N)$. \widetilde{D} is selfadjoint elliptic, so has an orthonormal basis of eigenfunctions φ_j with eigenvalues μ_j . Denote by $\overline{\varphi}_j$ the parallel translation of φ_j along the curves with tangents $\partial/\partial x$. Consider a metric of the form

$$(1.3) \quad dx^2 + h(x)^2 \widetilde{g}$$

where \widetilde{g} is the metric on N , and $h(x)$ a warping factor; in (0.1), $h(x) = (x^2 + t^2)^{1/2}$. According to [C, Prop. (2.5)], if $\widetilde{D}\varphi = \mu\varphi$, $\overline{\varphi}$ is the parallel

translation of φ , and f and g are functions of x alone, then

$$D \left(f\bar{\varphi} + g \frac{\partial}{\partial x} \cdot \bar{\varphi} \right) = \left(f' + \frac{n}{2} \frac{h'}{h} f - \frac{\mu}{h} g \right) \frac{\partial}{\partial x} \cdot \bar{\varphi} + \left(\frac{\mu}{h} f - g' - \frac{n}{2} \frac{h'}{h} g \right) \bar{\varphi}$$

where $n = \dim(N)$. Moreover, the norm of any section $f\bar{\varphi} + g\partial/\partial x \cdot \bar{\varphi}$ supported in $\{|x| < 1\}$ is

$$\|\varphi\|^2 \int_{-1}^1 (|f|^2 + |g|^2) h^n dx.$$

So if we introduce a normalizing factor $h^{-n/2}$ and represent a section as

$$(1.4) \quad s = \sum h^{-n/2} \left(f_j \bar{\varphi}_j + g_j \frac{\partial}{\partial x} \cdot \bar{\varphi}_j \right)$$

then

$$Ds = \sum h^{-n/2} \left[(h^{-1} \mu_j f_j - g_j') \bar{\varphi}_j + (f_j' - h^{-1} \mu_j g_j) \frac{\partial}{\partial x} \cdot \bar{\varphi}_j \right].$$

This gives the desired representation of D near the hypersurface N ; it is the direct sum of operators of the form

$$(1.5) \quad \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} \frac{\partial}{\partial x} + \frac{1}{h} \begin{bmatrix} \mu_j & 0 \\ 0 & -\mu_j \end{bmatrix} \\ = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} \left(\frac{\partial}{\partial x} + \frac{1}{h} \begin{bmatrix} 0 & -\mu_j \\ -\mu_j & 0 \end{bmatrix} \right)$$

acting on pairs $\begin{bmatrix} f \\ g \end{bmatrix}$ with norm $\int_{-1}^1 (|f(x)|^2 + |g(x)|^2) dx$. The pair $\begin{bmatrix} f_j \\ g_j \end{bmatrix}$ represents $h^{-n/2} (f_j \bar{\varphi}_j + g_j \partial/\partial x \cdot \bar{\varphi}_j)$, and $\tilde{D}\varphi_j = \mu_j \varphi_j$.

To represent $D^+ : C^\infty(S^+) \rightarrow C^\infty(S^-)$, we use the convenient fact that

$$\omega_{n+1} \cdot \left(\bar{\varphi} \pm \frac{\partial}{\partial x} \cdot \bar{\varphi} \right) = \pm \bar{\varphi} + \frac{\partial}{\partial x} \cdot \bar{\varphi}.$$

Thus S^\pm is spanned by sections $\{f_j(x)(\bar{\varphi}_j \pm \partial/\partial x \cdot \bar{\varphi}_j)\}$, and D^+ is a direct sum of operators (1.5) acting on pairs $\begin{bmatrix} f \\ g \end{bmatrix}$, which is equivalent to $\partial/\partial x - h^{-1} \mu_j$ acting on f .

In the limiting case, where $h(x) = |x|$, the analysis in [BS] shows that each eigenvalue μ_j in $(-1/2, 1/2)$ contributes a function f_j with a singularity x^{μ_j} as $x \rightarrow 0^+$, and $|x|^{-\mu_j}$ as $x \rightarrow 0^-$, yielding a section s_j in the maximal domain of D^+ with singularity $x^{-(n/2)+\mu_j}$ as $x \rightarrow 0^+$ and $|x|^{-(n/2)-\mu_j}$ as $x \rightarrow 0^-$. For the full Dirac operator in (1.5), the eigenvalues are $\pm \mu_j$, so there is a section on each side of N with singularity $|x|^{-(n/2) \pm \mu_j}$.

2. A parametrix for D_t near N . To construct a parametrix for the full Dirac operator D_t , we cancel the matrix $\begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}$ in (1.5), introduce eigenvectors $\begin{bmatrix} 1 \\ \pm 1 \end{bmatrix}$ for the matrix $\begin{bmatrix} 0 & -\mu \\ -\mu & 0 \end{bmatrix}$, and thus obtain operators of the form

$$(2.1) \quad \frac{\partial}{\partial x} \mp h^{-1}\mu.$$

So consider the equation

$$(2.2) \quad u'(x) - \mu h_t(x)^{-1}u(x) = f(x)$$

with f and u in L^2 , and $h_t(x) = (x^2 + t^2)^{1/2}$. The solution is

$$(2.3) \quad u(x) = (x + h_t(x))^\mu \left[\int_{?}^x (y + h_t(y))^{-\mu} f(y) dy + c \right]$$

where the questionable limit of integration is to be determined by our desire for a limit as $t \rightarrow 0$. Note that $x + h_t(x)$ is a positive and increasing function for real x . Thus in (2.3) we would like $x < y$ when $\mu > 0$, and $y < x$ when $\mu < 0$. So we define a parametrix for (2.2) on the interval $(-1, 1)$ by

$$(2.4) \quad Q_t f(x) = \begin{cases} (x + h_t(x))^\mu \int_1^x (y + h_t(y))^{-\mu} f(y) dy & \text{if } \mu > 0, \\ (x + h_t(x))^\mu \int_{-1}^x (y + h_t(y))^{-\mu} f(y) dy & \text{if } \mu \leq 0. \end{cases}$$

We will show that this has a limit as $t \rightarrow 0$. For the case $\mu > 0$ the limit is

$$(2.5) \quad Q_0 f(x) = \begin{cases} \int_1^x (x/y)^\mu f(y) dy, & x > 0, \\ \int_0^x |y/x|^\mu f(y) dy, & x < 0, \end{cases} \quad \mu > 0.$$

For $\mu < 0$ the lower limits are 0 for $x > 0$ and -1 for $x < 0$; for $\mu = 0$, we have merely $Q_0 f(x) = \int_{-1}^x f(y) dy$. (With $\mu = 0$, we could just as well integrate from 1 as from -1 .)

LEMMA 1. For $t < 1$,

$$(2.6) \quad \|Q_t\| \leq C(1 + |\mu|)^{-1},$$

$$(2.7) \quad \|Q_t - Q_0\| \leq C(1 + |\mu|)^{-1}(t + t^\mu) \log(1 + 1/t),$$

with C independent of μ and t .

Proof of (2.6). The case $\mu = 0$ is clear, since then Q_t is independent of t ; and the operator Q_t for $\mu < 0$ is minus the adjoint of the one for $-\mu$; so it suffices to take $\mu > 0$. For this case, we split both domain and range as the direct sum of $L^2(-1, 0)$ and $L^2(0, 1)$.

For $0 \leq x \leq y$ we have

$$(2.8) \quad \frac{x + h_t(x)}{y + h_t(y)} \leq \frac{x + t}{y + t}.$$

So the integral operator on $L^2(0, 1)$ with kernel

$$(x + h_t(x))^\mu (y + h_t(y))^{-\mu}, \quad 0 \leq x \leq y \leq 1,$$

has norm bounded by

$$(2.9) \quad \sup_{x \leq 1} (x + t)^\mu \int_x^1 (y + t)^{-\mu} dy + \sup_{y \leq 1} (y + t)^{-\mu} \int_0^y (x + t)^\mu dx \leq \frac{C}{1 + \mu}$$

with C independent of μ and t , $0 \leq t \leq 1$. Likewise, the operator from $L^2(0, 1)$ to $L^2(-1, 0)$ defined by

$$Kf(x) = (x + h_t(x))^\mu \int_0^1 (y + h_t(y))^{-\mu} f(y) dy, \quad -1 \leq x \leq 0,$$

has norm

$$(2.10) \quad \begin{aligned} \|K\| &\leq \left[\int_{-1}^0 (x + h_t(x))^{2\mu} dx \int_0^1 (y + h_t(y))^{-2\mu} dy \right]^{1/2} \\ &= t^{2\mu} \int_0^1 (y + h_t(y))^{-2\mu} dy \leq t^{2\mu} \int_0^1 (y + t)^{-2\mu} dy \\ &= \frac{t}{2\mu - 1} \left[1 - \left(\frac{t}{1 + t} \right)^{2\mu - 1} \right]. \end{aligned}$$

(In the integral from -1 to 0 , replace x by $-y$.) When $\mu \geq 1$ this is $\leq Ct/(1 + \mu)$. When $\mu < 1$, the Mean Value Theorem applied to $(\frac{t}{1+t})^x$ gives $\|K\| \leq t \log(1 + 1/t)$. Thus, for $0 \leq t \leq 1$,

$$(2.11) \quad \|K\| \leq \frac{C}{1 + \mu} t \log(1 + 1/t).$$

This gives the desired estimate for Q_t from $L^2(0, 1)$ to $L^2(-1, 1)$. The part of Q_t acting from $L^2(-1, 0)$ to $L^2(0, 1)$ is 0, and the part from $L^2(-1, 0)$ to $L^2(-1, 0)$ is equivalent, by an obvious change of variable, to an operator on $L^2(0, 1)$ with kernel

$$(x + h_t(x))^{-\mu} (y + h_t(y))^\mu, \quad 0 \leq y \leq x \leq 1.$$

This is the adjoint of the one estimated in (2.9), so has the same norm.

Proof of (2.7). As before, it suffices to take $\mu > 0$. In view of (2.11), and the adjoint relation exploited at the end of the previous proof, it suffices to consider the operator on $L^2(0, 1)$ with kernel

$$\left[\frac{x + h_t(x)}{y + h_t(y)} \right]^\mu - \left[\frac{x}{y} \right]^\mu, \quad 0 \leq x \leq y \leq 1.$$

We must show that this operator satisfies the estimate (2.7). For brevity, we drop the subscript t on h_t . Using (2.8) again, it suffices to treat the kernel

$$k(x, y) = \left(\frac{x+t}{y+t}\right)^\mu - \left(\frac{x}{y}\right)^\mu, \quad 0 \leq x \leq y \leq 1.$$

By Schur's Test, the norm of this operator is dominated by

$$\begin{aligned} & \sup_{x \leq 1} \int_x^1 k(x, y) dy + \sup_{y \leq 1} \int_0^y k(x, y) dx \\ &= \frac{1}{1-\mu} [t^\mu(1+t)^{1-\mu} - t] + \frac{t}{\mu+1} \left[1 - \left(\frac{t}{1+t}\right)^\mu \right], \quad \mu \neq 1. \end{aligned}$$

The second term is $\leq t/(\mu+1)$. The first is $\leq t/(1-\mu)$ for $\mu > 1$, and the Mean Value Theorem shows that it is $\leq t^{\min(\mu,1)} \log(1+1/t)$. This proves (2.7), and with it Lemma 1.

3. The global parametrix, and graph continuity. This section is essentially a review of the corresponding part of [SS], *mutatis mutandis*.

Since D_t is selfadjoint, $i + D_t$ is invertible. We construct the resolvent $(i + D_t)^{-1}$ and show that it converges in norm as $t \rightarrow 0$, thus proving graph continuity of the family D_t at $t = 0$.

Let Q_i be a parametrix for the Dirac operator D_t on the part of M where $|x| > 1/2$, the "interior" of M . Since the metric (0.1) varies with t , so will Q_i vary with t ; but it is easy to guarantee that Q_i varies smoothly in t , for $t \geq 0$. Thus

$$(3.1) \quad D_t Q_i = I + T_i$$

where the kernel of T_i is C^∞ on $M \times M \times [0, 1]$, and the kernel of Q_i is C^∞ off the diagonal. The parametrices Q_t in (2.4), acting in the various eigenspaces, can be combined to give a parametrix for D_t in $\{|x| < 1\}$; we denote this, too, by Q_t . Each operator in (2.4) is compact, and the direct sum of these operators converges in norm, because of the factor $(1 + \mu)^{-1}$ in (2.6), so the combined operator Q_t is also compact.

We now set

$$(3.2) \quad P_t = \varphi_i Q_i \psi_i + \varphi_c Q_t \psi_c$$

where $\psi_i + \psi_c = 1$, $\varphi_i \psi_i = \psi_i$, $\varphi_c \psi_c = \psi_c$, while φ_i and ψ_i vanish for $|x| < 1/2$, and φ_c and ψ_c are supported in $\{|x| < 1\}$. We have

$$(3.3) \quad (i + D_t) P_t = I + R_t$$

where the compact remainder is

$$R_t = iP_t + \varphi_i T_i \psi_i + \varphi'_i e_0 \cdot Q_i \psi_i + \varphi'_c e_0 \cdot Q_t \psi_c$$

with e_0 denoting Clifford multiplication by $\partial/\partial x$. Since $I + R_t$ is Fredholm with index zero, and $i + D_t$ is invertible, it follows that P_t is a Fredholm operator, with index zero, from $L^2(S)$ to the domain of D_t . Using the representation (1.4), we can think of P_t as acting in a fixed Hilbert space, independent of t . Define D_0 to be the Dirac operator D in the conic limiting metric $g_0 = dx^2 + x^2\tilde{g}$, with domain defined by

$$(3.4) \quad \lim_{x \rightarrow 0^+} |x|^{n/2} s(x) = \lim_{x \rightarrow 0^-} |x|^{n/2} s(x),$$

both limits existing (as in (0.2)). One checks that both P_0 and P_0^* map into this domain; in eigenspaces with $\mu \neq 0$ both limits are 0, while in the 0 eigenspace they are equal. Moreover, (3.4) defines a selfadjoint realization of D , and (3.3) holds true for $t = 0$.

Let N_0 be the nullspace of P_0 , and $k = \dim(N_0)$. Since P_0 has index zero as a map into $\text{Dom}(D_0)$, there is a k -dimensional subspace, V , of $\text{Dom}(D_0)$, which is linearly independent of $\text{Range}(P_0)$. We will show in Lemma 3 below that V can be taken as a subspace of

$$C_c^\infty(S) = \text{sections of } S \text{ vanishing in a neighborhood of } N.$$

(This is the argument omitted from [SS].) Let V_0 map N_0 isomorphically onto V , and N_0^\perp to 0. Then $P_0 + V_0$ is an isomorphism of $L^2(S)$ onto $\text{Dom}(D_0)$, so

$$(i + D_t)(P_t + V_0) = I + R_t + (i + D_t)V_0$$

is invertible when $t = 0$. Since V_0 is in $C_c^\infty(S)$, the right hand side varies continuously with t , and

$$(3.5) \quad (i + D_t)^{-1} = (P_t + V_0)(I + R_t + (i + D_t)V_0)^{-1}$$

varies continuously with t , as $t \rightarrow 0$. This proves Theorem 1.

Now suppose that $\dim(M)$ is even, so there is a direct sum $S = S^+ \oplus S^-$, with D_t mapping $C^\infty(S^\pm)$ to $C^\infty(S^\mp)$. Denote by D_t^+ the restriction of D_t to the sections of S^+ . We can represent the full Dirac operator D_t as $\begin{bmatrix} 0 & D_t^+ \\ D_t^- & 0 \end{bmatrix}$, with $D_t^- = (D_t^+)^*$. Hence the continuity of $(i + D_t)^{-1}$ implies

that of $\begin{bmatrix} I & iD_t^+ \\ -(iD_t^+)^* & I \end{bmatrix}^{-1}$, and hence the graph continuity of D_t^+ (see [SS]). This proves Theorem 2, and with it Corollary 1.

Now consider the situation of Corollary 2 of Theorem 2, where N divides M into two disjoint parts; denote the resulting conic limit manifolds by \overline{M}_+ and \overline{M}_- . Denote by $\nu(\tilde{D})$ the nullity of the Dirac operator \tilde{D} on N . Then the condition (3.4) defining the domain of the conic limiting operator involves $\nu(\tilde{D})$ matching conditions. These can be replaced by an equal

number of vanishing conditions, by requiring

$$\lim_{x \rightarrow 0^+} |x|^{n/2} s(x) = 0, \quad \lim_{x \rightarrow 0^-} |x|^{n/2} s(x) \text{ exists.}$$

Since the number of conditions is the same, the index is the same. Moreover, these conditions define precisely the Dirichlet realization on \overline{M}_+ and the Neumann realization on \overline{M}_- . Hence

$$\begin{aligned} \widehat{A}(M) &= \text{ind}(D_0) = \text{ind}(D_{\text{Dir}}^+(\overline{M}_+)) + \text{ind}(D_{\text{Neu}}^+(\overline{M}_-)) \\ &= \widehat{A}_{\text{Dir}}(\overline{M}_+) + \widehat{A}_{\text{Neu}}(\overline{M}_-). \end{aligned}$$

If we replace the condition “ $\lim_{x \rightarrow 0^-} |x|^{n/2} s(x)$ exists” by the condition “ $\lim_{x \rightarrow 0^-} |x|^{n/2} s(x) = 0$ ”, we add $\nu(\tilde{D})$ boundary conditions, and thus decrease the index by $\nu(\tilde{D})$. Thus we obtain the final statement of Corollary 2,

$$\widehat{A}(M) = \widehat{A}_{\text{Dir}}(\overline{M}_+) + \widehat{A}_{\text{Dir}}(\overline{M}_-) + \nu(\tilde{D}).$$

It remains to prove that the complementary space V above can be taken in $C_c^\infty(S)$. This depends on two lemmas.

LEMMA 2. *If u is in $\text{Dom}(D_0)$ and orthogonal (in the graph norm for D_0) to both $\text{Range}(P_0)$ and $C_c^\infty(S)$, then $u = 0$.*

Proof. The conditions on u are

$$(3.6) \quad (u, v) + (D_0 u, D_0 v) = 0 \quad \text{for all } v \text{ in } C_c^\infty(S),$$

$$(3.7) \quad (u, P_0 f) + (D_0 u, D_0 P_0 f) = 0 \quad \text{for all } f \text{ in } L^2(S).$$

The first condition shows that $u + D_0^2 u = 0$ except perhaps on N , so $D_0 u$ is in the maximal domain of D . We will show that $D_0 u$ is actually in the domain of D_0 , so the equation $u + D_0^2 u = 0$ implies that $u = 0$.

For f supported in $\{|x| < 1/2\}$, $D_0 P_0 f = f$, so the second condition (3.7) gives

$$(3.8) \quad D_0 u = -P_0^* u \quad \text{in } \{|x| < 1/2\}.$$

But (like P_0) P_0^* maps into the domain of D_0 . Since the domain of D_0 is defined by a limiting condition as $x \rightarrow 0^\pm$, it follows from (3.8) that $D_0 u$ is in the domain of D_0 , and this completes the proof of Lemma 2.

LEMMA 3. *$\text{Range}(P_0)$ has a linear complement spanned by elements of $C_c^\infty(S)$.*

Proof. By Lemma 2, $\text{Range}(P_0) + C_c^\infty(S)$ is dense in $\text{Dom}(D_0)$. But $\text{Range}(P_0)$ has finite codimension in that space, so $\text{Range}(P_0) + C_c^\infty(S)$ is also a closed subspace, and thus gives all of $\text{Dom}(D_0)$. Now if $\{u_1, \dots, u_k\}$ span a linear complement of $\text{Range}(P_0)$ in $\text{Dom}(D_0)$, then each $u_j = P_0 f_j + v_j$ for a v_j in $C_c^\infty(S)$; then the v_j span a linear complement of $\text{Range}(P_0)$.

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Reçu par la Rédaction le 25.4.1990